



## Interface and Self Assembly

*Nick Fang*

Course Website: [nanoHUB.org](http://nanoHUB.org)  
[Compass.illinois.edu](http://Compass.illinois.edu)



# Proposal (Oct 12-16)



- Please send me your intended topic by tomorrow
- Based on similarity of topics you may be asked to work/present as group
- A tentative schedule will be posted on Friday

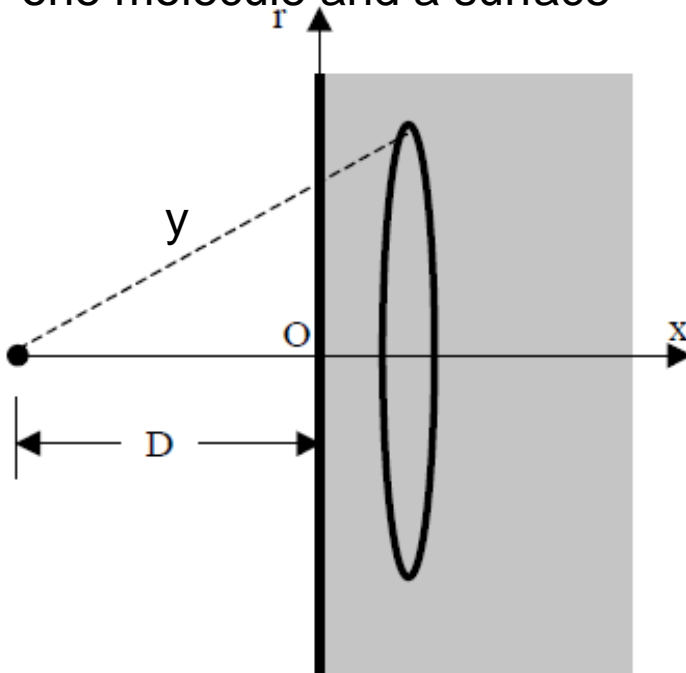


# Origin of Adhesion: VdW Forces



- Van de Waals force: the long range interactions between molecules
- Recall the potential energy:  $\phi = -\frac{\alpha}{r^6} + \frac{\beta}{r^n}$  ( $n \approx 12$ )

Let's find the interaction between one molecule and a surface



At distance  $x$  in the wall, consider a circle of radius  $r$ :

$$y = \sqrt{(D+x)^2 + r^2}$$

Number density of molecules in wall

$$\Phi(D) = \int_0^{\infty} \int_0^{\infty} \rho \phi(y) 2\pi r dr dx$$

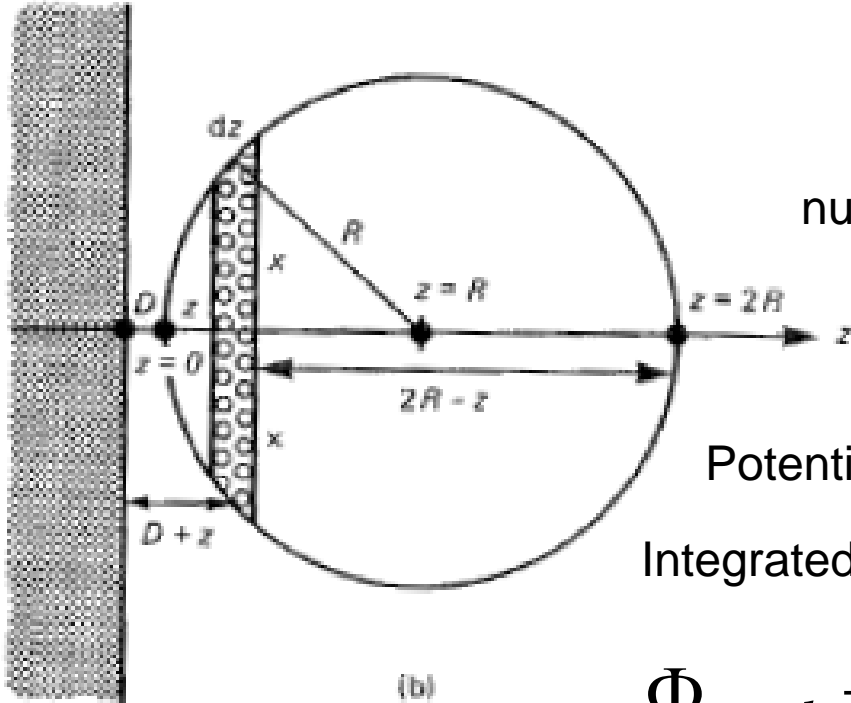
$$\Phi(D) = -\frac{\pi\rho\alpha}{6D^3}$$



# Example: Sphere to Flat Surface



Consider a thin sheet at location  $z$  on the sphere:



$$\text{Radius} \quad x = \sqrt{(2R - z)z}$$

number of molecules on this sheet:

$$\rho_2 \pi x^2 dz = \rho_2 \pi (2R - z) z dz$$

$$\text{Potential energy per molecule: } \Phi(D+z) = -\frac{\pi \rho_1 \alpha}{6(D+z)^3}$$

Integrated over whole sphere:

$$\Phi_{total} = -\rho_1 \rho_2 \pi^2 \alpha \int_0^{2R} \frac{(2R - z)z}{6(D+z)^3} dz$$

When  $D \ll R$ ,  $2R - z \sim 2R$

$$\Phi_{total}(D) \approx -\frac{\rho_1 \rho_2 \alpha \pi^2 R}{6D}$$

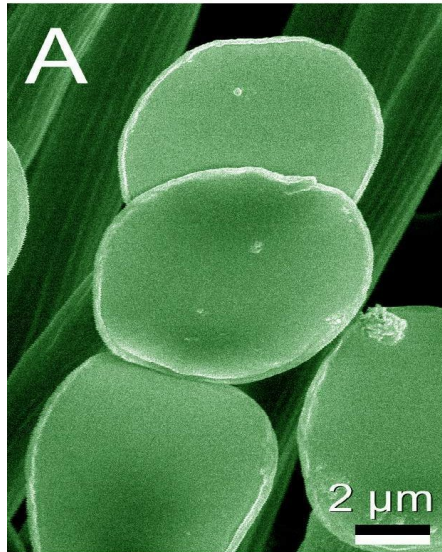
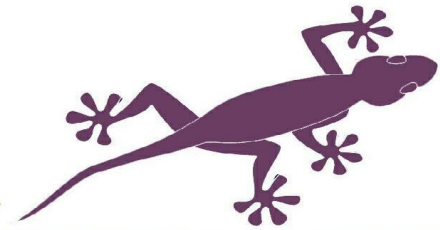
← Size of contact  
← Distance of contact



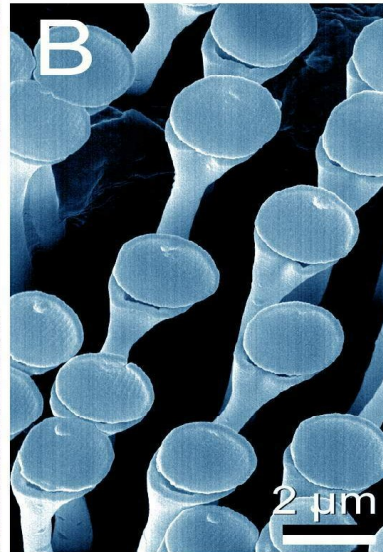
# Adhesion Enhancement by Nano-toes



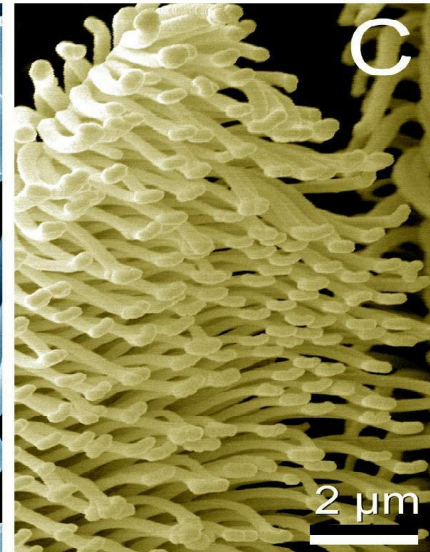
body mass →



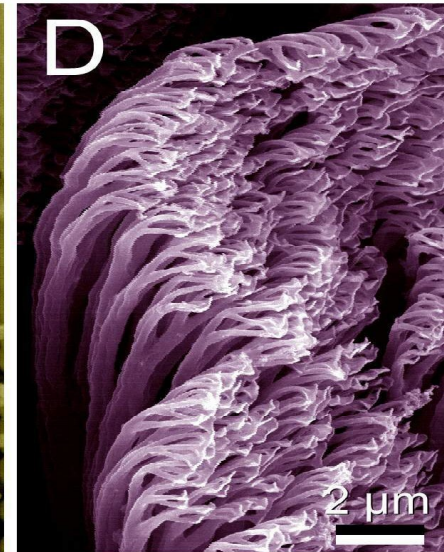
beetle



fly



spider



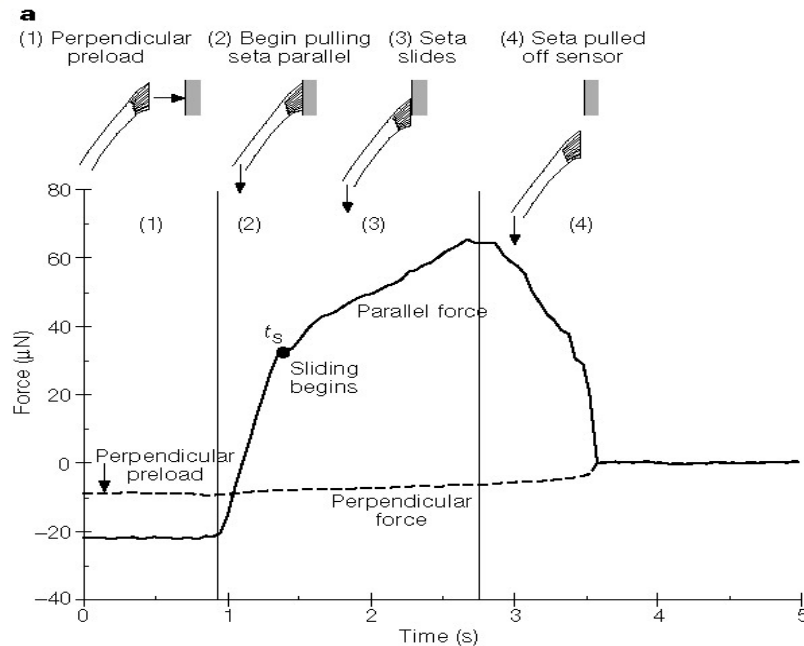
gecko

<http://shasta.mpi-stuttgart.mpg.de/research/Bio-tribology.htm>

$$\Phi_{total}(D) \approx -\frac{\rho_1 \rho_2 \alpha \pi^2 R}{6D}$$

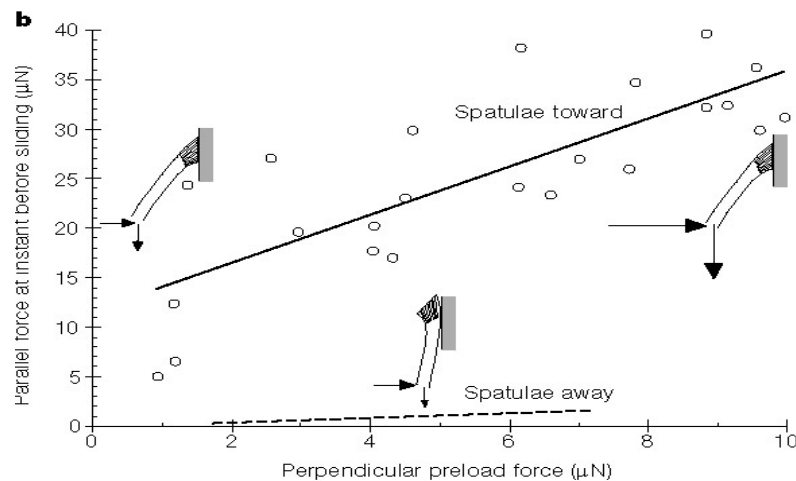


# How Can a Gecko Lift Its Foot Off?



**“These lizards uncurl their toes like a paper party favor whistle when putting their feet down;  
- and peel the toes back up as if removing a piece of tape when they step away.”**

Chemical & Engineering News, 2000

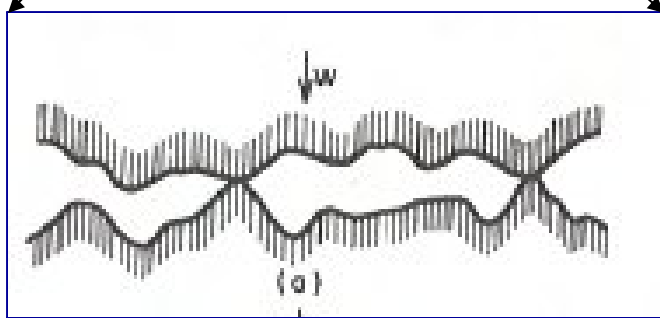
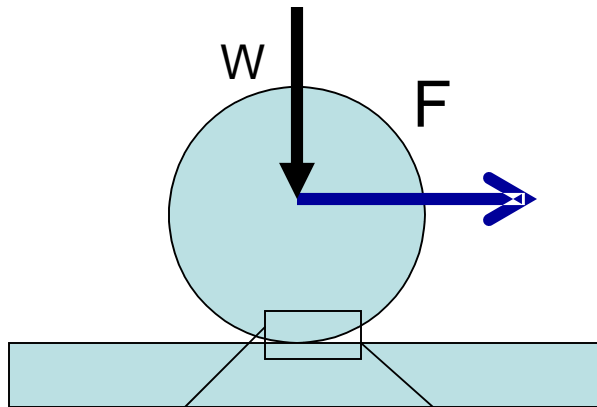


K. Autumn et al, Nature 405, 681-685(2000)





# Microscopic View of Friction



$$F \approx \frac{A}{D} (\underbrace{\gamma_A}_{\text{Surface energy in Advancing contact}} - \underbrace{\gamma_R}_{\text{Surface energy in Receding contact}})$$

Surface energy  
in Advancing  
contact

Surface energy  
in Receding  
contact

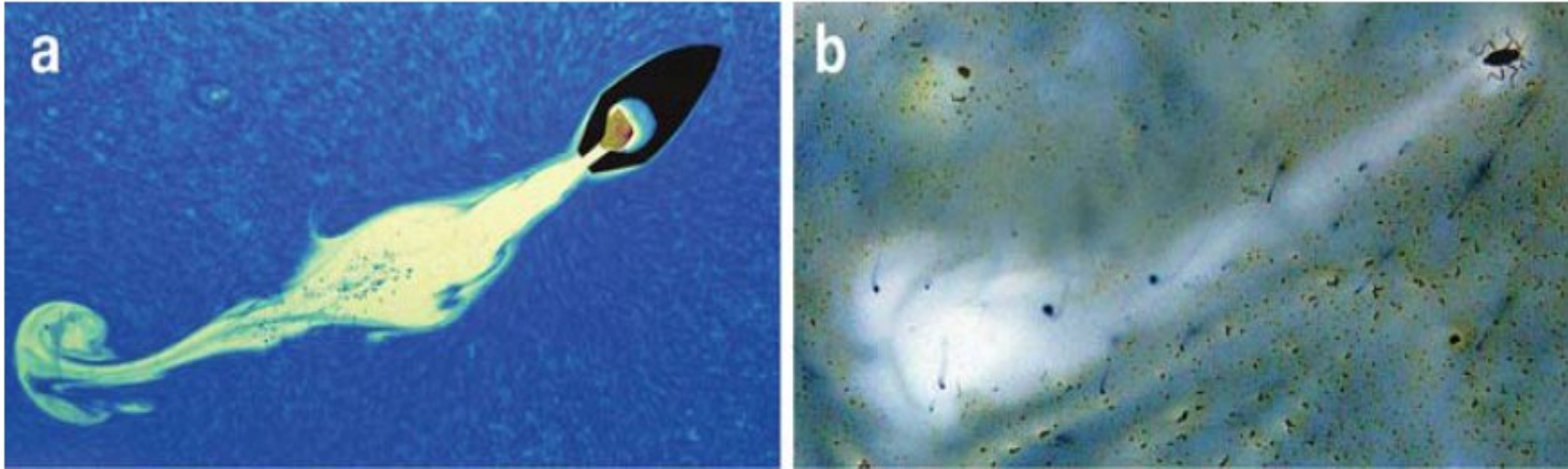
**Derjaguin (1957) proposed  
correction of friction**

$$F \approx \mu W + \underline{\mu A p_0}$$

Due to adhesion  
energy (no external  
force needed)



# From Surface Forces to Self Assembly



From Bush, J. W. M., and D. L. Hu. “WALKING ON WATER: Biocomotion at the Interface.” *Annu. Rev. Fluid Mech.* 2006. 38:339–69

“By releasing a surfactant, water striders and other insects was able to propel itself toward and up the meniscus”





# Outline



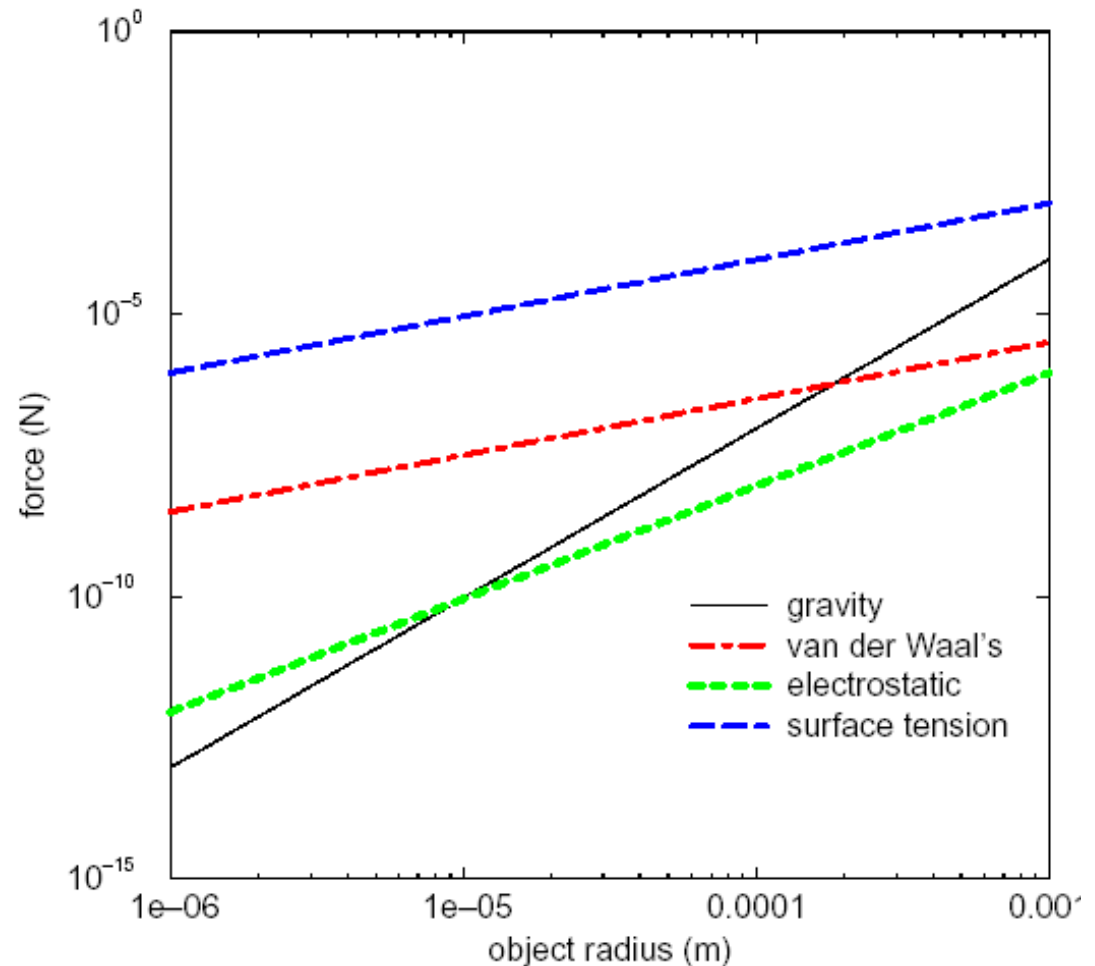
- **Self Assembly**
  - Thermodynamics of Micelle self assembly
  - Micro/Meso/Macroscale Self Assembly
  - Limitations



# Driving Forces for Self-Assembly

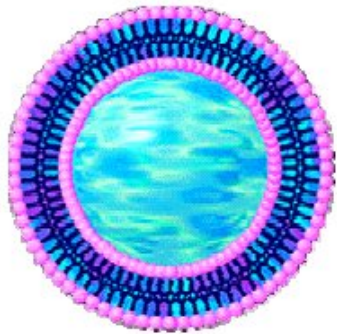


- Molecular Bonding Forces
- Steric Energy
- Capillary Forces
- Electrostatic Forces
- Magnetic Forces

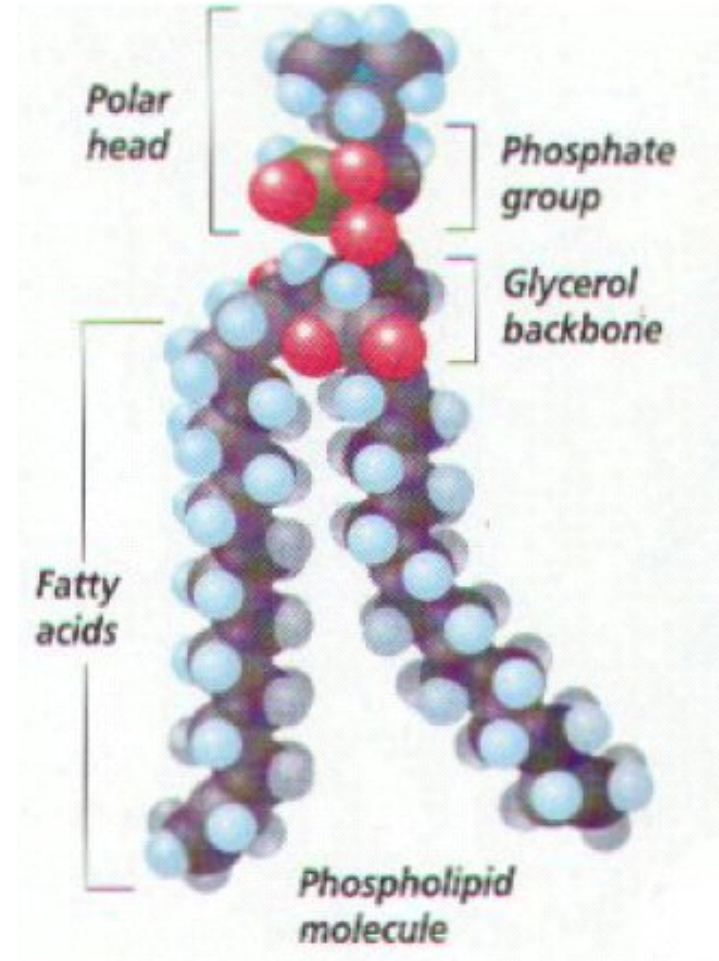
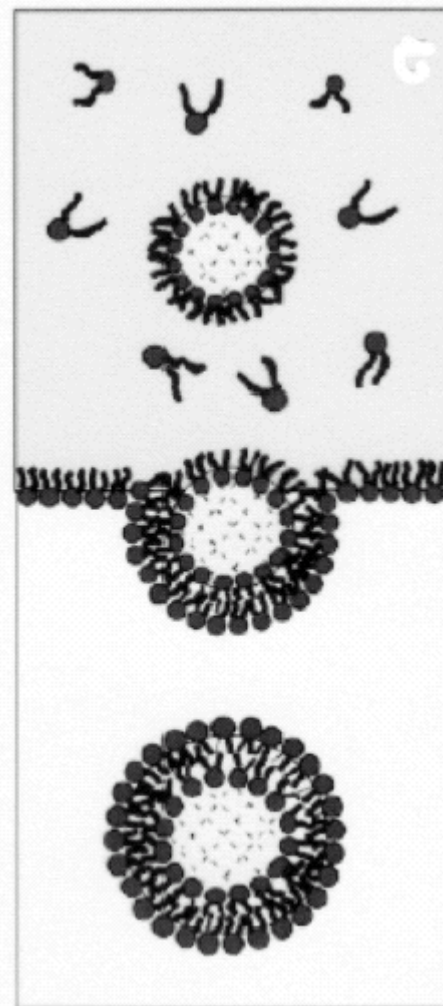




# Assembly by Surface Energy



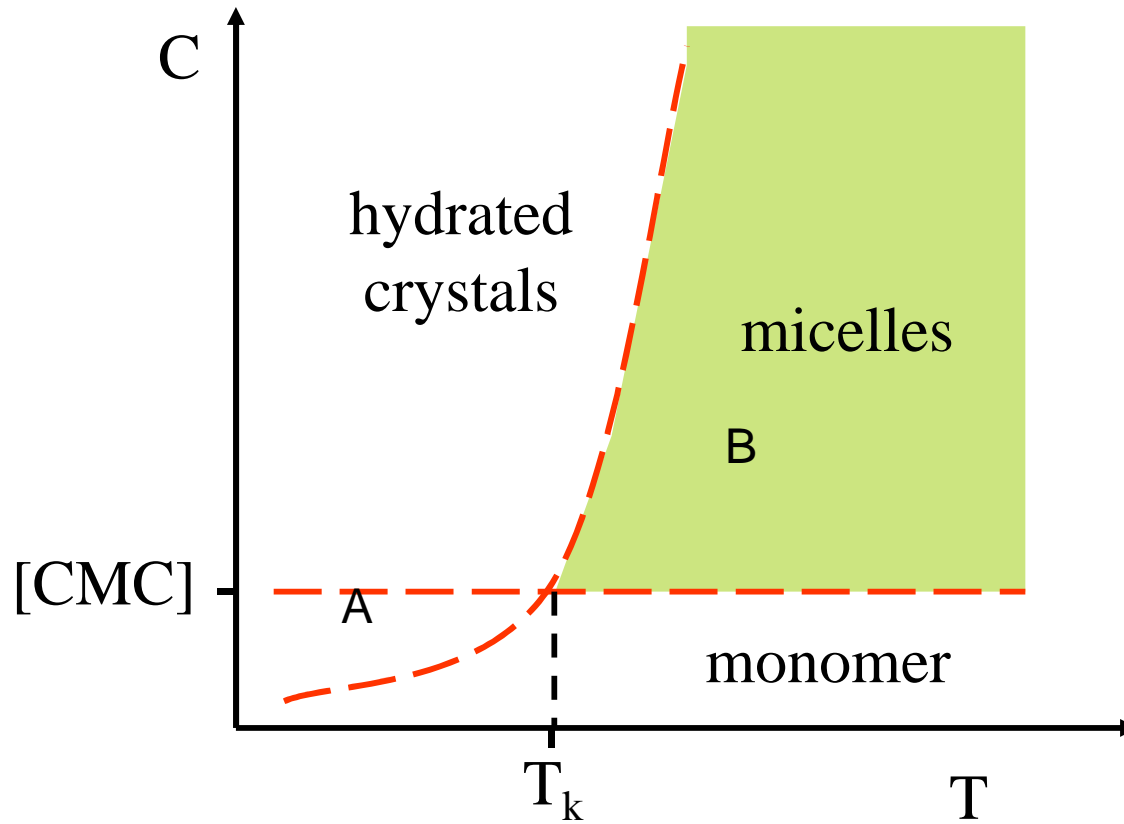
**Vesicles of bilayer membranes**



**Surfactant molecule: amphiphilic**

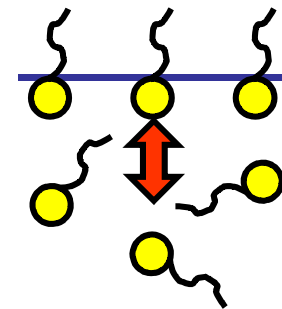


# Phase Diagram of Micelles

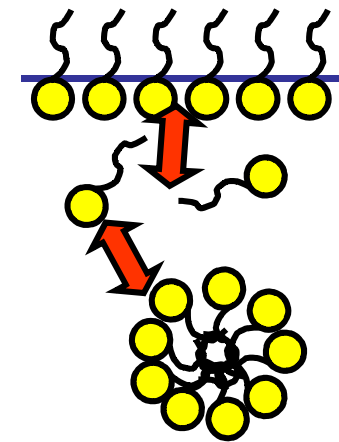


Phase diagram

Situation A



Situation B

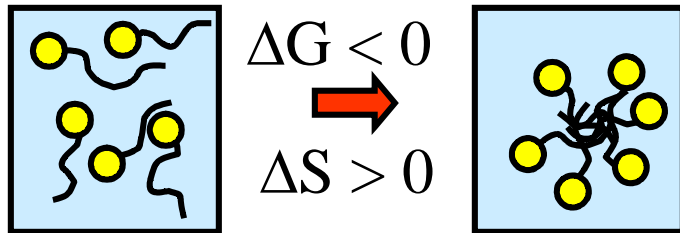




# Micelle Properties



## 1. Hydrophobic effect

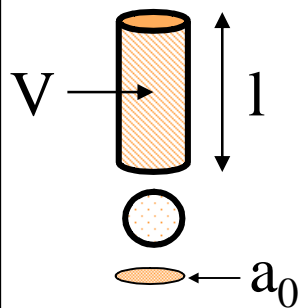


## 2. Head repulsion



Electrostatic or Steric

## 3. Packing efficacy (geometric factor)



$$N_s = \frac{V}{a_0 l}$$

$$N_s = 0.33$$

$$N_s = 0.5$$

$$N_s = 1.0$$



bilayer



# Molecular Self Assembly Geometries



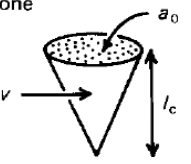


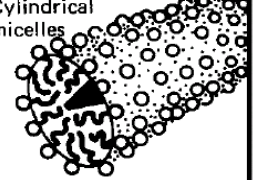
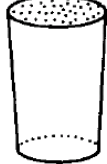
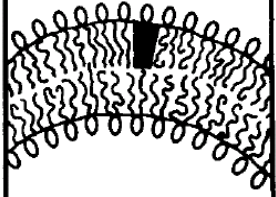
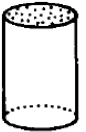
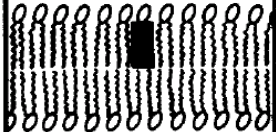
Surfactant molecular shape/interactions mainly determines aggregate geometry.

Critical packing factor =  $v/a_0 l_c$  (unitless), where:

$v$  = molecular volume of surfactant chain

$a_0$  = area per surfactant head

$l_c$  = length of surfactant chain

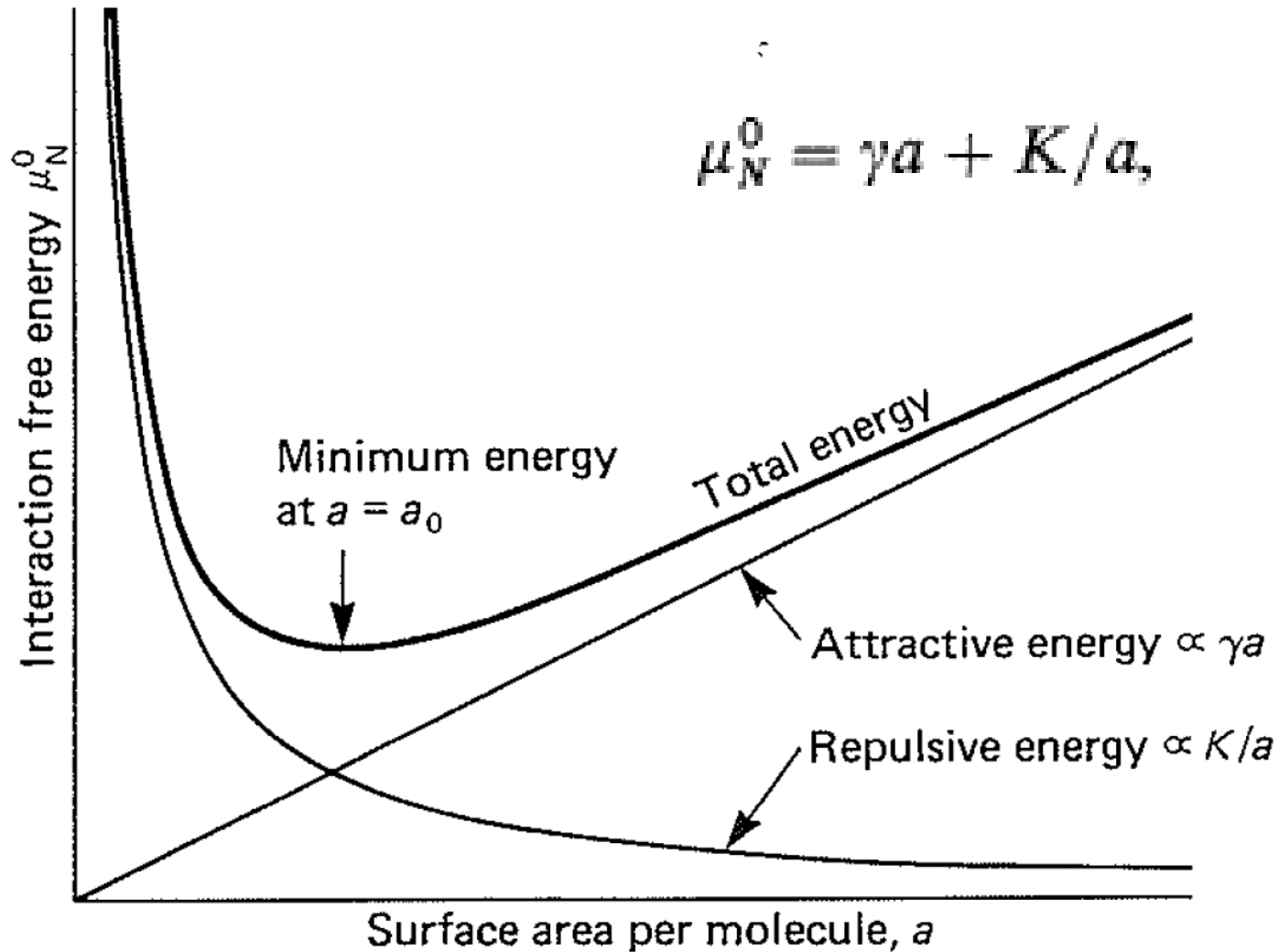
Lipid	Critical packing parameter $v/a_0 l_c$	Critical packing shape	Structures formed
Single-chained lipids (surfactants) with large head-group areas: <i>SDS in low salt</i>	$< 1/3$	Cone 	Spherical micelles 
Single-chained lipids with small head-group areas: <i>SDS and CTAB in high salt, nonionic lipids</i>	$1/3-1/2$	Truncated cone 	Cylindrical micelles 
Double-chained lipids with large head-group areas, fluid chains: <i>Phosphatidyl choline (lecithin), phosphatidyl serine, phosphatidyl glycerol, phosphatidyl inositol, phosphatidic acid, sphingomyelin, DGDG*, dihexadecyl phosphate, dialkyl dimethyl ammonium salts</i>	$1/2-1$	Truncated cone 	Flexible bilayers, vesicles 
Double-chained lipids with small head-group areas, anionic lipids in high salt, saturated frozen chains: <i>phosphatidyl ethanolamine, phosphatidyl serine + Ca<sup>2+</sup></i>	$\sim 1$	Cylinder 	Planar bilayers 

From Israelachvili, Chap 16





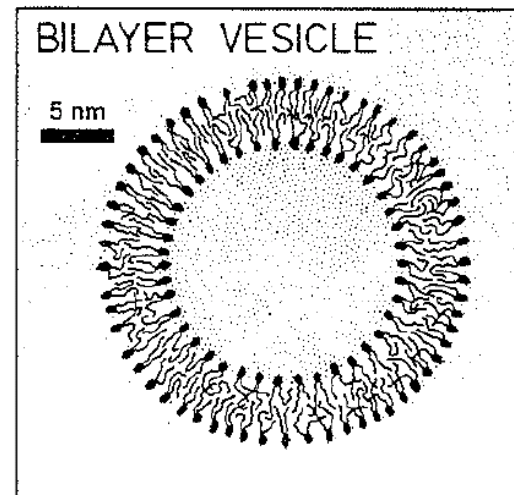
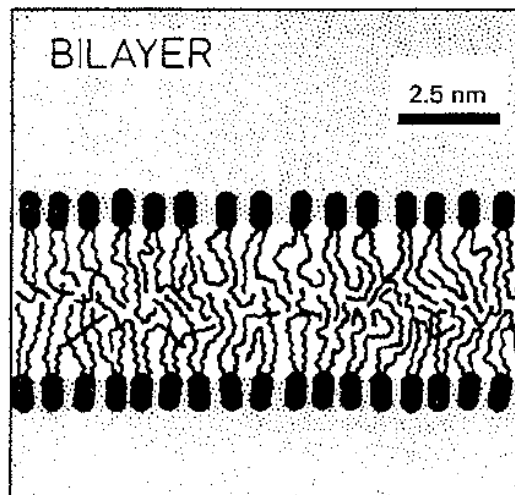
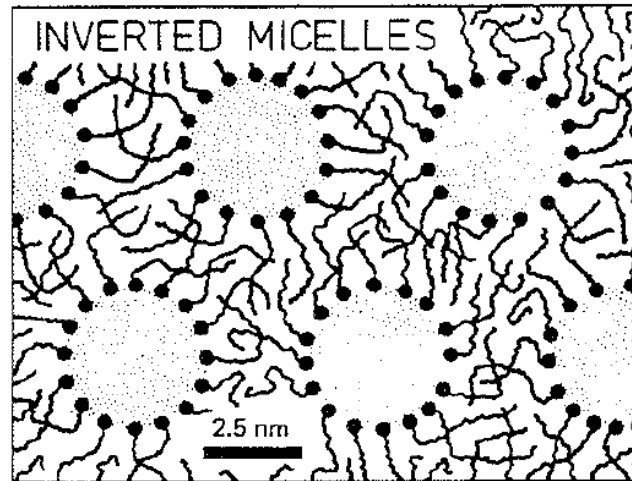
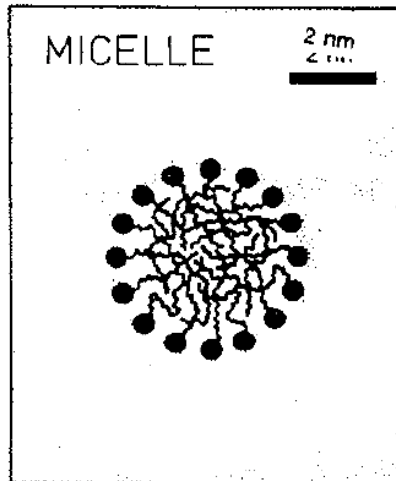
# Critical Size of Self Assembled Micelles



From Israelachvili, Chap 17



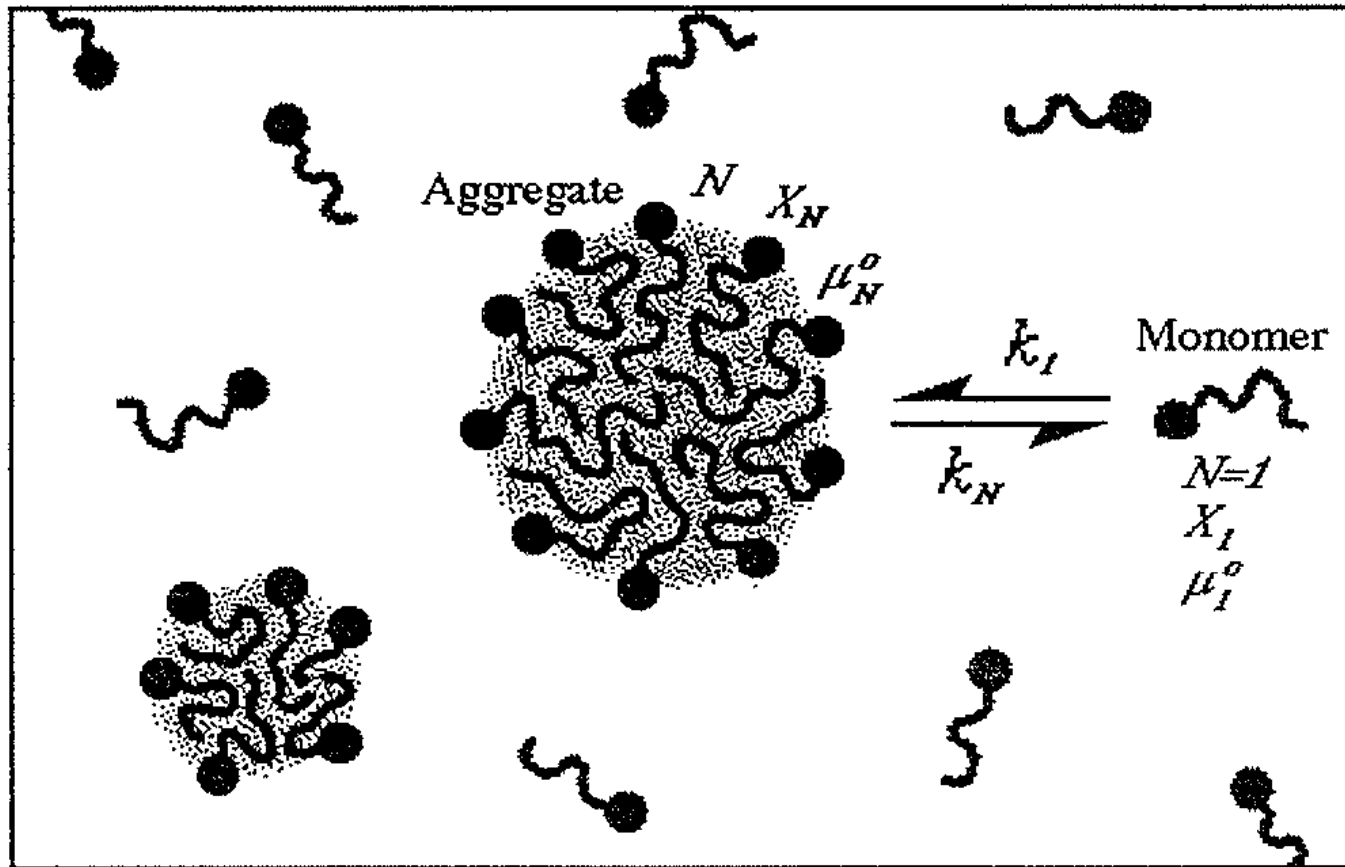
# Possible Phase Transformation of SAMs



From Israelachvili, Chap 16



# Thermodynamic Balance of SAM

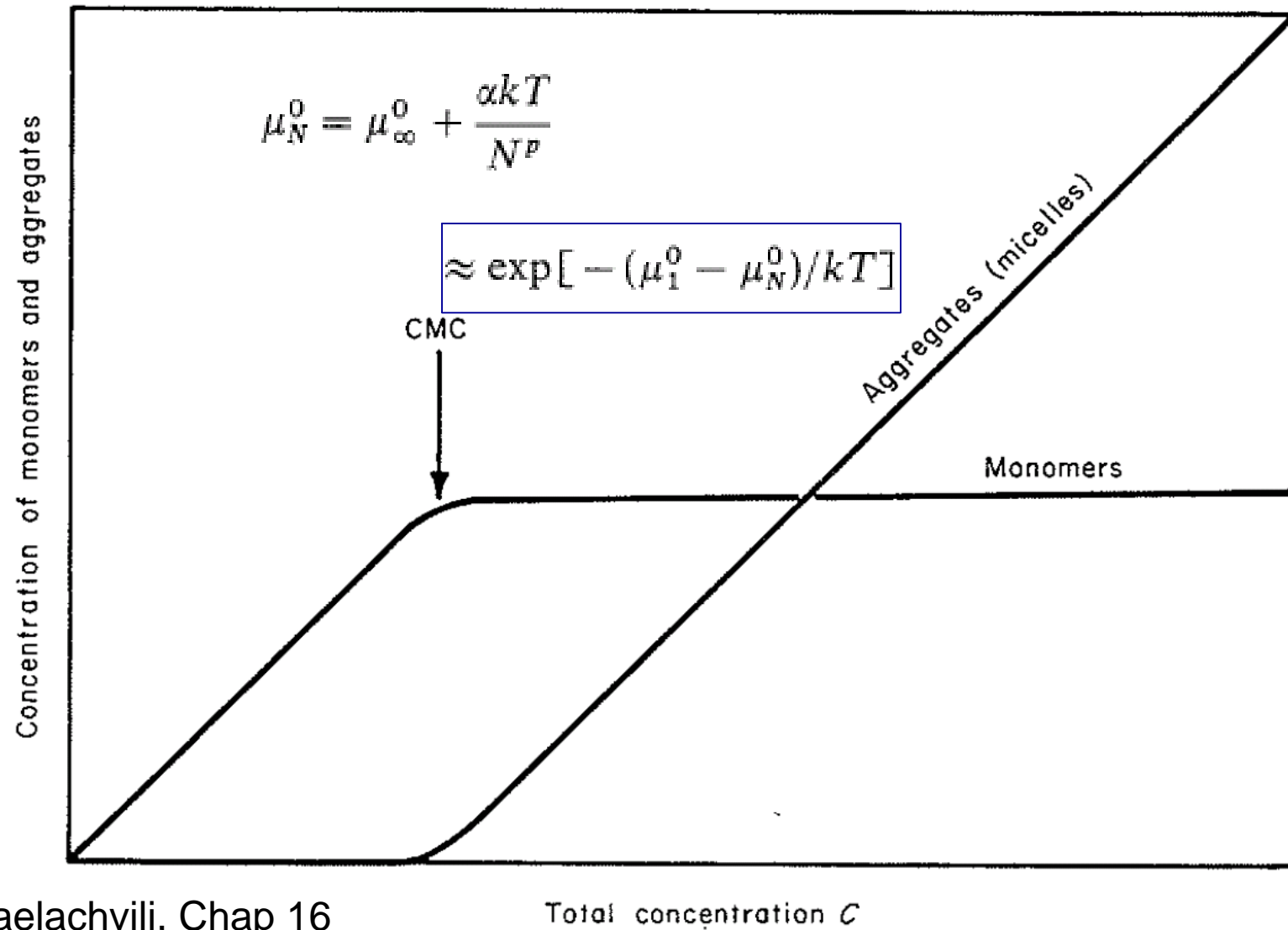


**Fig. 16.3.** Association of  $N$  monomers into an aggregate (e.g., a micelle). The mean lifetime of an amphiphilic molecule in a small micelle is very short, typically  $10^{-5}$ – $10^{-3}$  s.

From Israelachvili, Chap 16



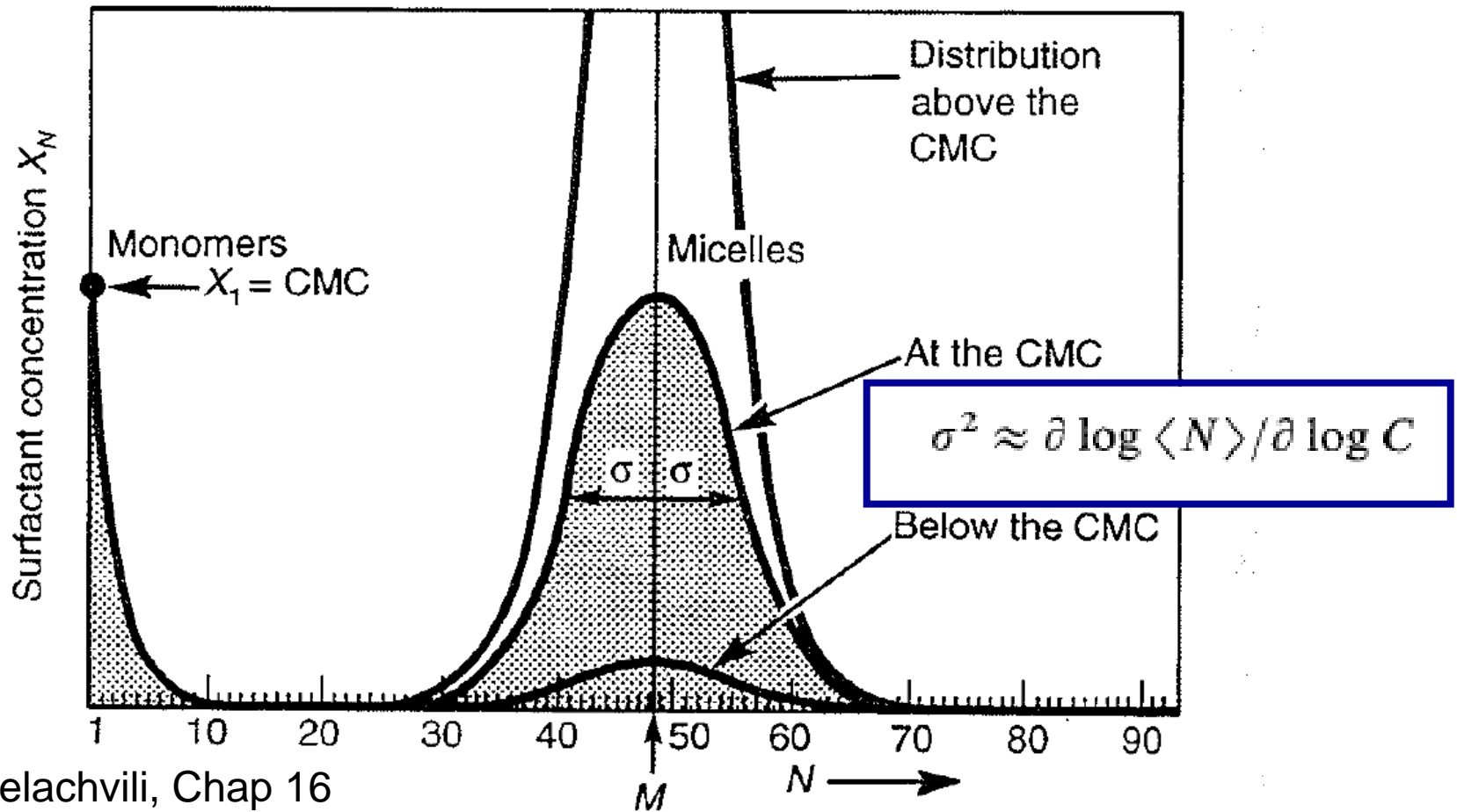
# Critical Micelle Concentration (CMC)



From Israelachvili, Chap 16



# Size Distribution of SAMs



From Israelachvili, Chap 16

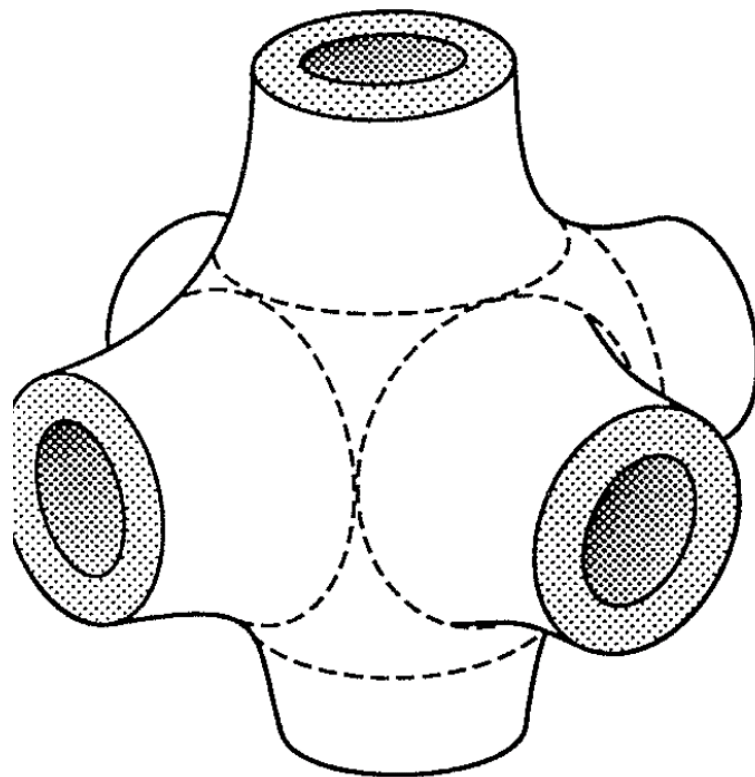
***Size deviation scales as square root of number ( $M$ ) of monomers required to form the micelle***



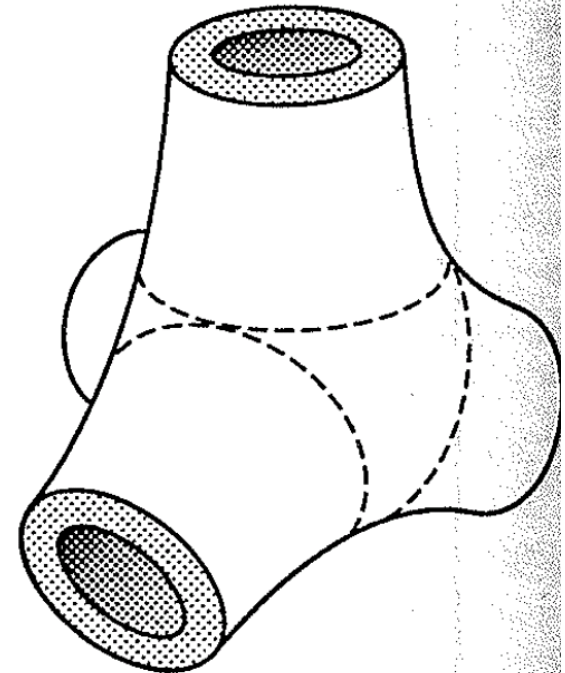
# More Complex Shapes



- Need to consider the curvature elasticity of membranes/micropHases



**A**



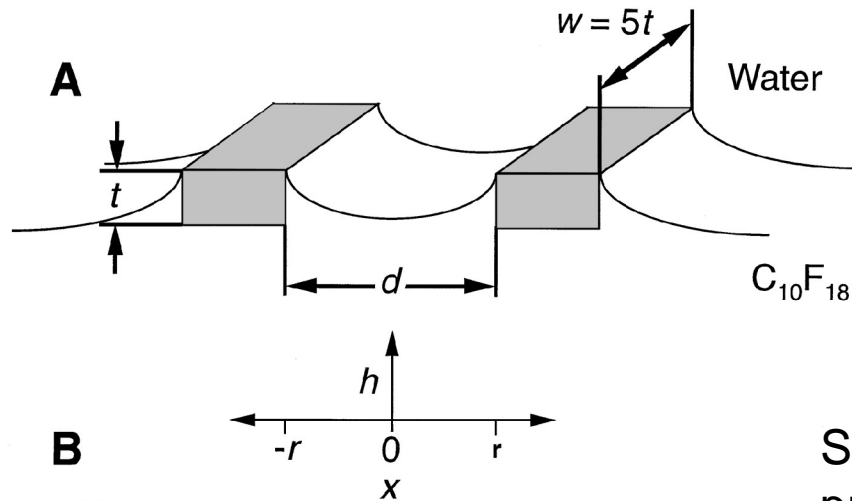
**B**

From Israelachvili, Chap 17





# Self Assembly Driven by Capillary Forces

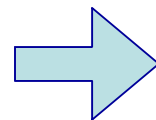
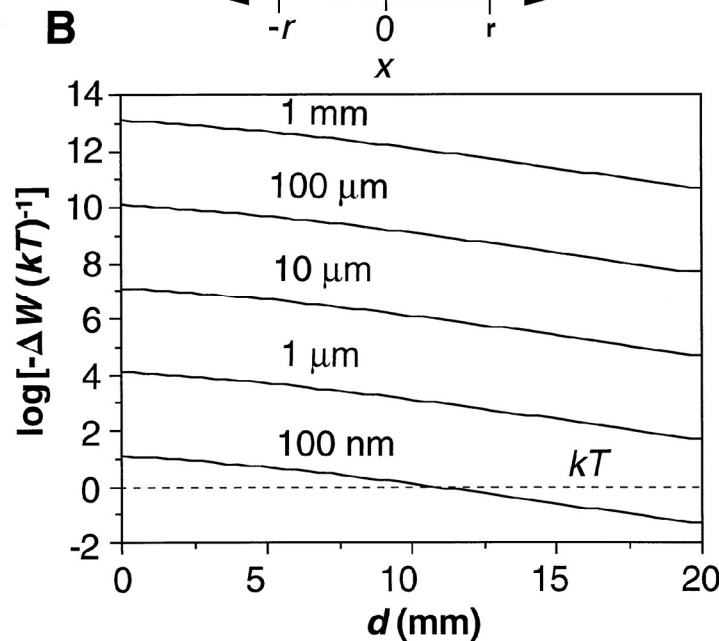


The height of the meniscus is given by Laplace Equation:

$$\frac{\partial^2 h}{\partial x^2} = \frac{1}{\gamma} (\Delta \rho g h - \Delta P_0)$$

Solution of the above gives the surface profile due to capillary force and gravity:

$$h(x) = t \left[ \frac{2}{1 - e^{(d/x_c)}} + \frac{e^{(-x/x_c)} + e^{(x/x_c)}}{e^{(d/2x_c)} - e^{(-d/2x_c)}} \right]$$



*self-assembly are favorable for objects with  $t$  as small as 100 nm*

WHITESIDES et al, Science 1997



# Additional Readings



- Jacob N. Israelachvili, “Intermolecular and Surface Forces”, Chapter 16, 17, Academic Press, 2nd Edition, 1992
- MRS Bulletin, Focused Issue on “Self Assembly in Materials Synthesis”, 2005
- Whitesides Group Website:  
<http://gmwgroup.harvard.edu/research.html>