



# **ECE606: Solid State Devices**

## **Lecture 16: Carrier Transport**

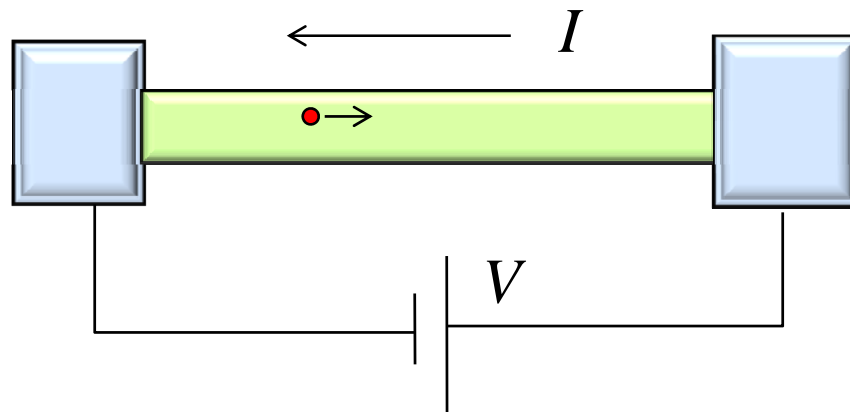
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# Outline

- 1) **Overview**
- 2) Drift Current
- 3) Physics of Mobility
- 4) High field effects
- 5) Conclusion

REF: Advanced Device Fundamentals, Pages 175- 192

# Current Flow Through Semiconductors



$$I = G \times V$$

$$= q \times n \times v \times A$$

Carrier  
Density

velocity

Depends on chemical composition,  
crystal structure, temperature, doping, etc.

## Quantum Mechanics + Equilibrium Statistical Mechanics

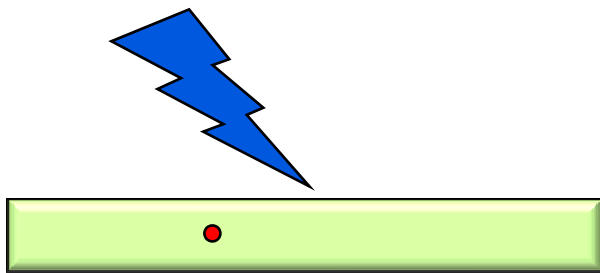
⇒ Encapsulated into concepts of effective masses  
and occupation factors (Ch. 1-4)

## Transport with scattering, non-equilibrium Statistical Mechanics

⇒ Encapsulated into drift-diffusion equation with  
recombination-generation (Ch. 5 & 6)

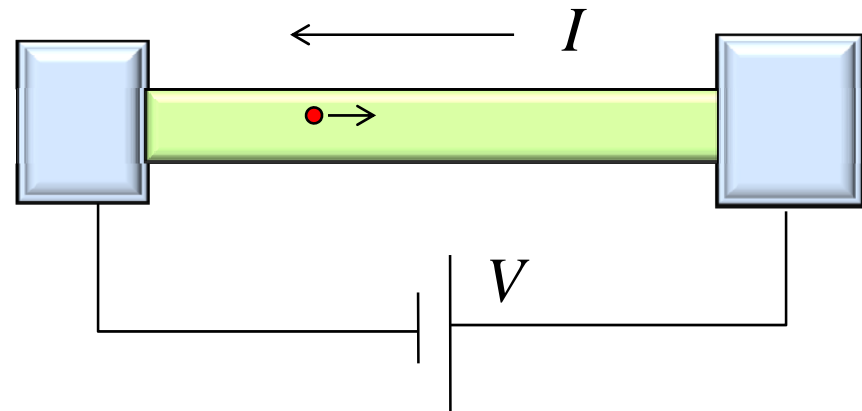
# Non-equilibrium Systems

Chapter 5



vs.

Chapter 6



# Summary of Transport Equations ...

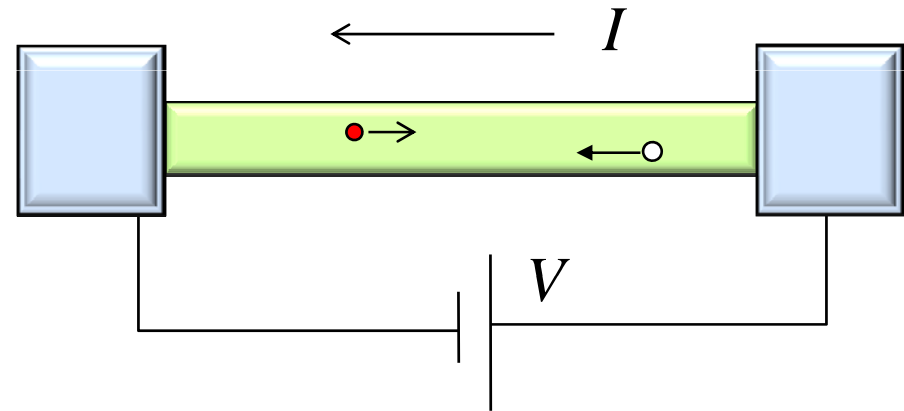
$$\nabla \cdot \mathbf{D} = q(p - n + N_D^+ - N_A^-)$$

$$\frac{\partial n}{\partial t} = \frac{1}{q} \nabla \cdot \mathbf{J}_N - r_N + g_N$$

$$\mathbf{J}_N = qn\mu_N \mathbf{E} + qD_N \nabla n$$

$$\frac{\partial p}{\partial t} = -\frac{1}{q} \nabla \cdot \mathbf{J}_P - r_P + g_P$$

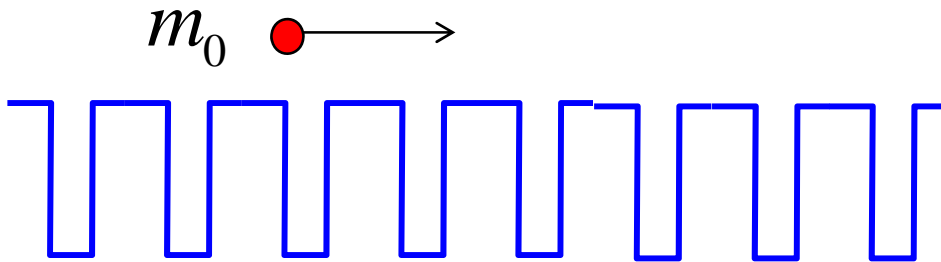
$$\mathbf{J}_P = qp\mu_P \mathbf{E} - qD_P \nabla p$$



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# Meaning of Effective Mass ...



$$\left( -\frac{\hbar^2}{2m_0} \frac{d^2}{dx^2} + U_{\text{crys}}(x) + U_{\text{ext}}(x) \right) \psi = E\psi$$



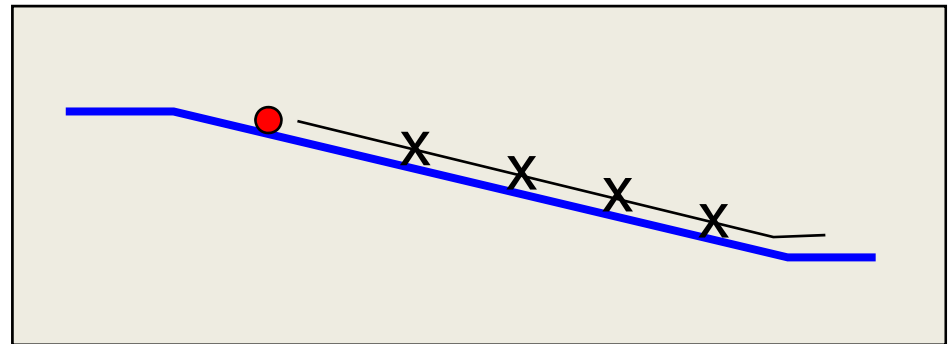
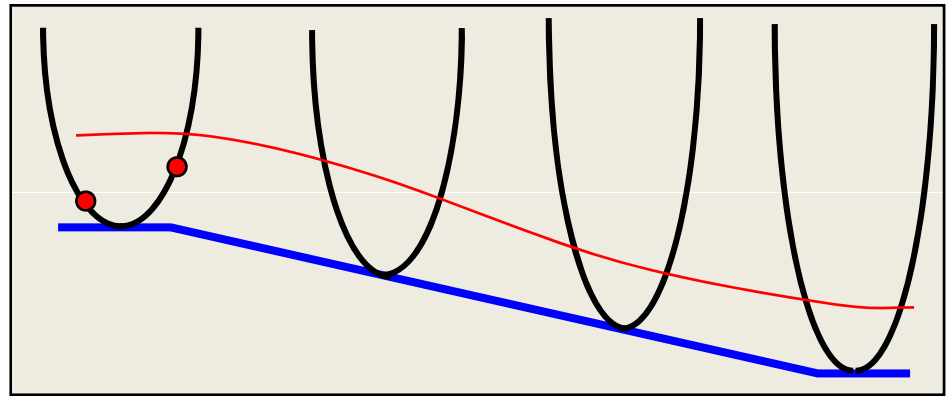
$$\left( -\frac{\hbar^2}{2m_n^*} \frac{d^2}{dx^2} + U_{\text{ext}}(x) \right) \phi = E\phi$$

# Drift by Electric field ....

$$J_n = qn\mu_n \mathcal{E}$$

$$\frac{d(m_n^* v)}{dt} = -q\mathcal{E} - \frac{m_n^* v}{\tau_n}$$

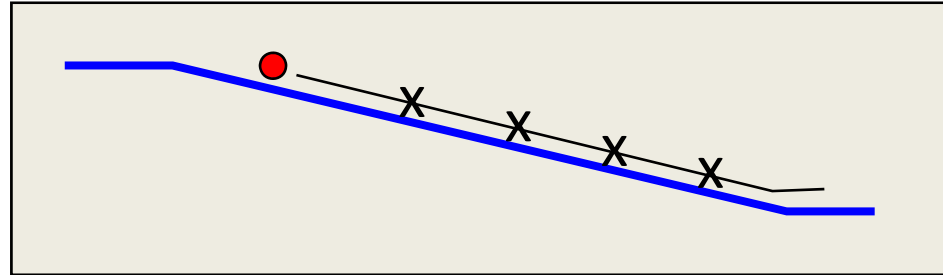
$$v(t) = -\frac{q\tau_n}{m_n^*} \mathcal{E} \left[ 1 - e^{-\frac{t}{\tau_n}} \right]$$





# Drift by Electric field ...

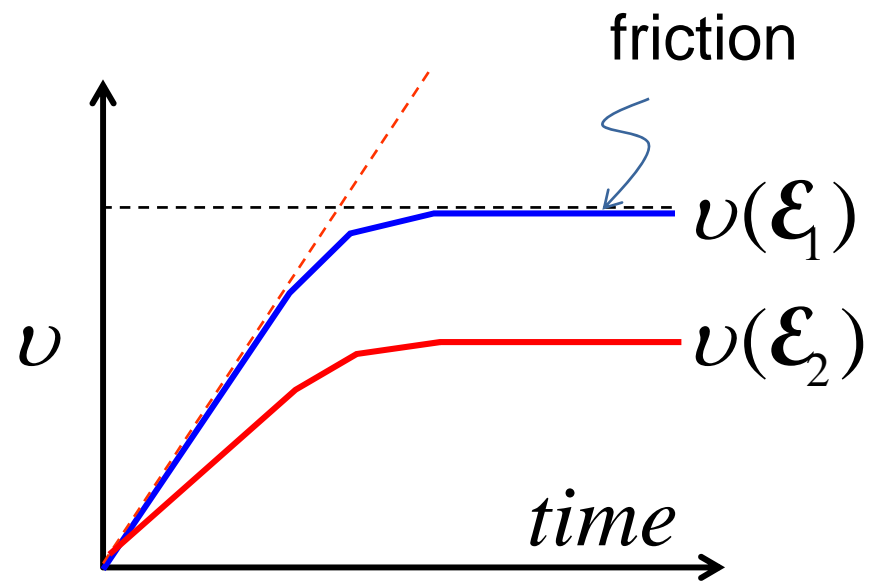
$$v(t) = -\frac{q\tau_n}{m_n^*} \mathcal{E} \left[ 1 - e^{-\frac{t}{\tau_n}} \right]$$



$$= -\frac{q\tau_n}{m_n^*} \mathcal{E} \quad (t \rightarrow \infty, 1-2 \text{ ps})$$

$$\equiv \mu_n \mathcal{E}$$

$$J_n = qn\mu_n \mathcal{E}$$

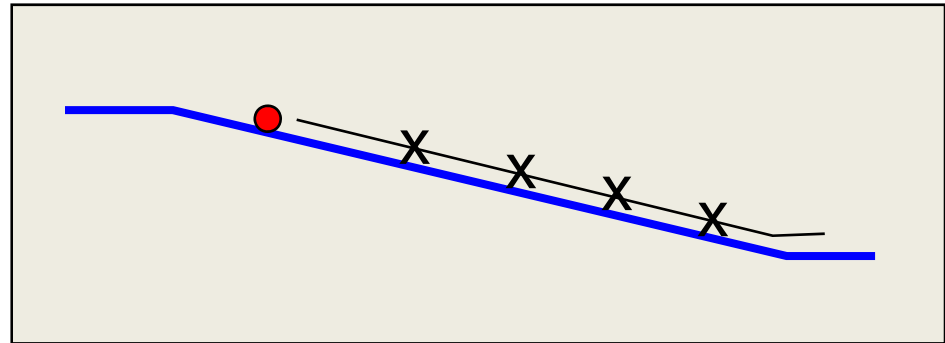


(Theory valid once  $t > 1-2 \text{ ps}$ )

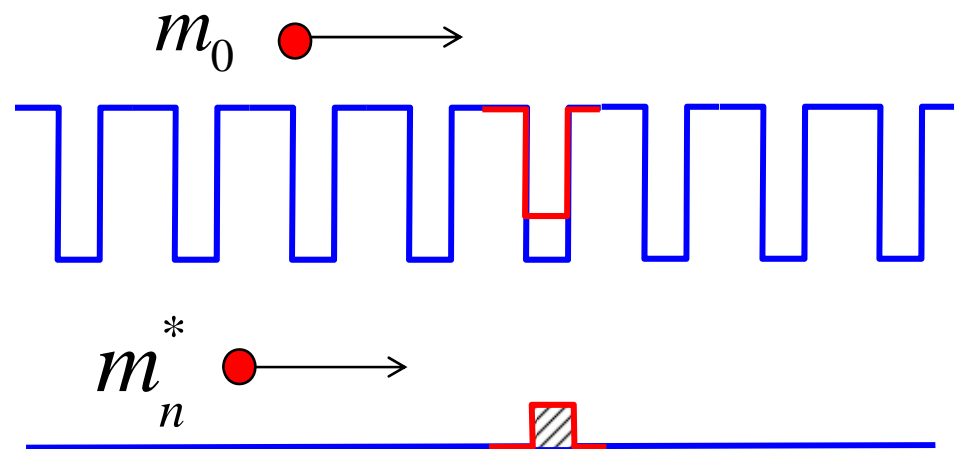
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# Mobility and Physics of Scattering Time



$$\mu_n = \frac{q\tau_n}{m_n^*}$$



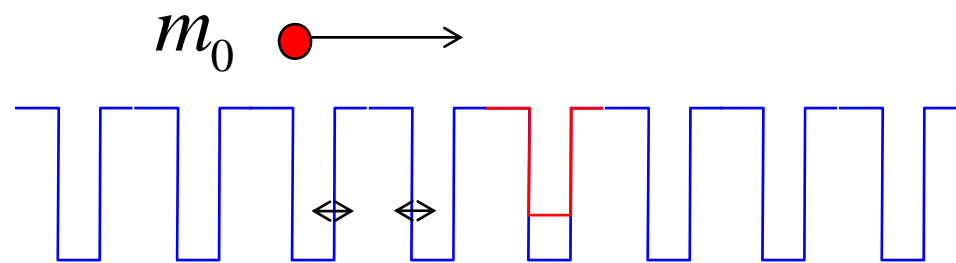
Fermi's Golden rule ...

$$\tau_n^{-1} \sim \left| \frac{2\pi}{\hbar} \int_{-\infty}^{\infty} \psi^*(x) U(x) \psi(x) dx \right|^2$$

# Phonon and Ionized Impurity Scattering

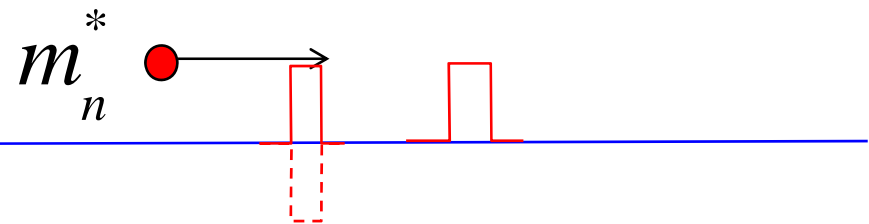
Ionized impurity

$$\tau_n \sim \frac{T^{3/2}}{N_D}$$



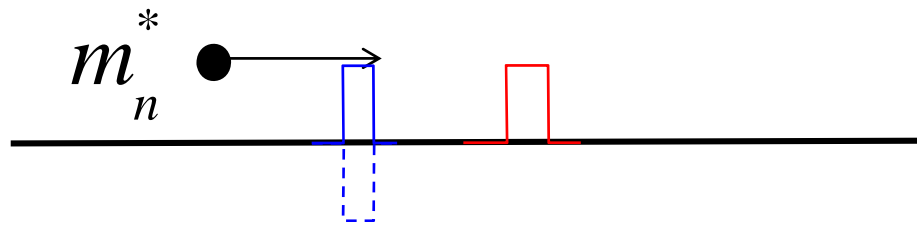
Higher temperature,  
more phonon scattering

$$\tau_n \sim T^{-3/2}$$



# Multiple Scattering Events

- Ionized impurity
- Phonon scattering
- others ....



$$\frac{1}{\mu_n} = \frac{1}{\mu_{ph}} + \frac{1}{\mu_{II}}$$

$$\Rightarrow \mu_n = \frac{\mu_{ph}\mu_{II}}{\mu_{ph} + \mu_{II}}$$

$$= \mu_{\min} + \left( \frac{\mu_{ph}\mu_{II}}{\mu_{ph} + \mu_{II}} - \mu_{\min} \right)$$

$$= \mu_{\min} + \left( \frac{\mu_0}{1 + (N_I/N_0)^\alpha} \right)$$

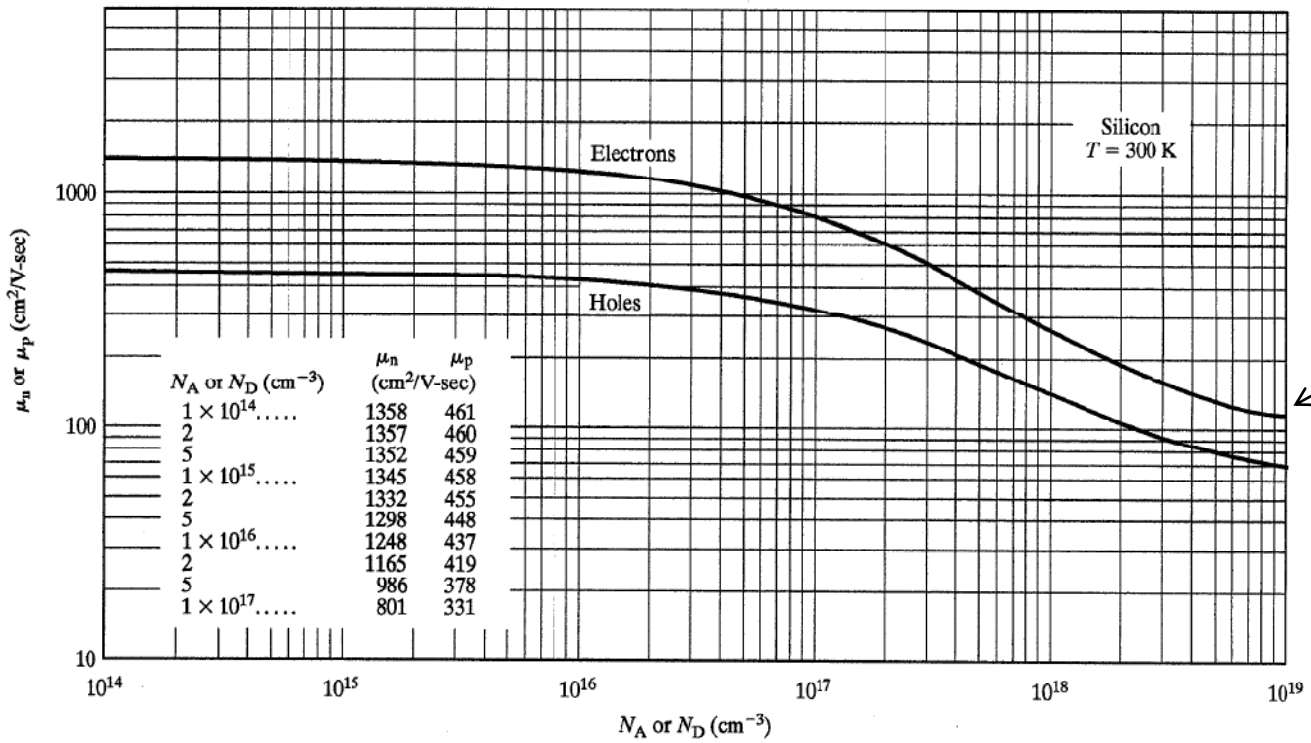
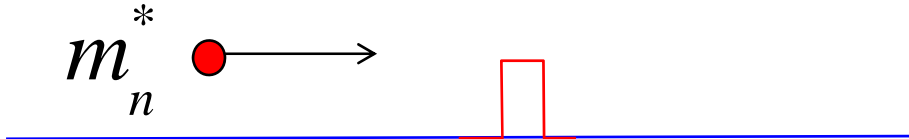
$$\frac{1}{\tau_n} = \frac{1}{\tau_{II}} + \frac{1}{\tau_{ph}} + \frac{1}{\tau_s} + \dots$$

$$\frac{1}{\mu_n} = \frac{m_n^*}{q\tau_n}$$

Matthesson Rule ....

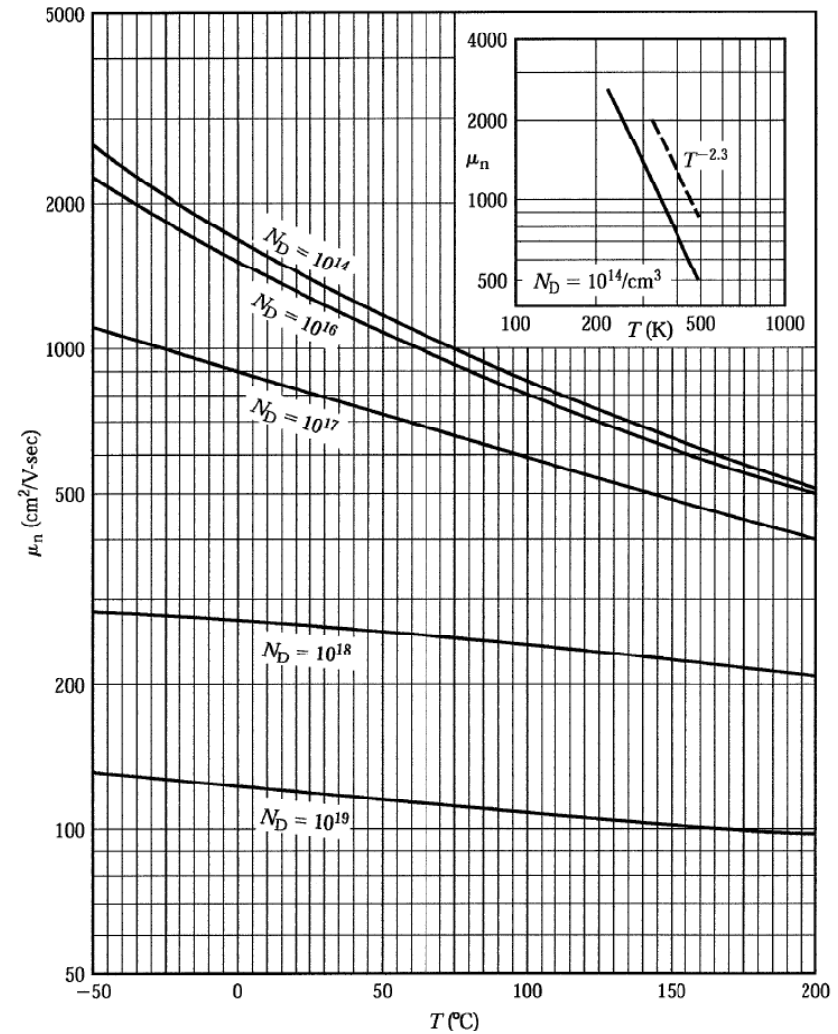
# Model for Ionized impurity Scattering

$$\mu_n = \mu_{n,\min} + \left( \frac{\mu_{0,n}}{1 + (N_I/N_{0,n})^{\alpha_n}} \right)$$



# Temperature-dependent Mobility

$$\mu_n \sim \tau_n \sim T^{-3/2}$$

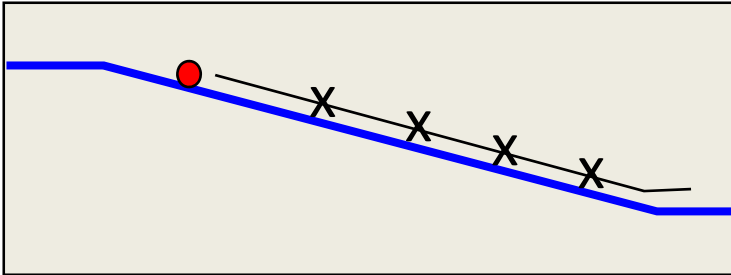


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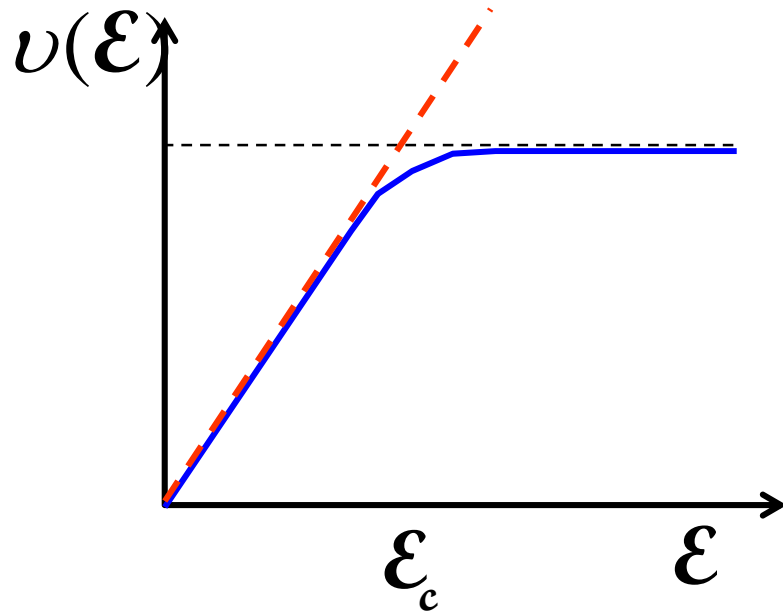
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# Mobility at High Fields?



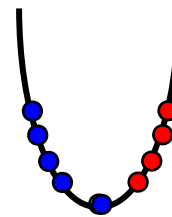
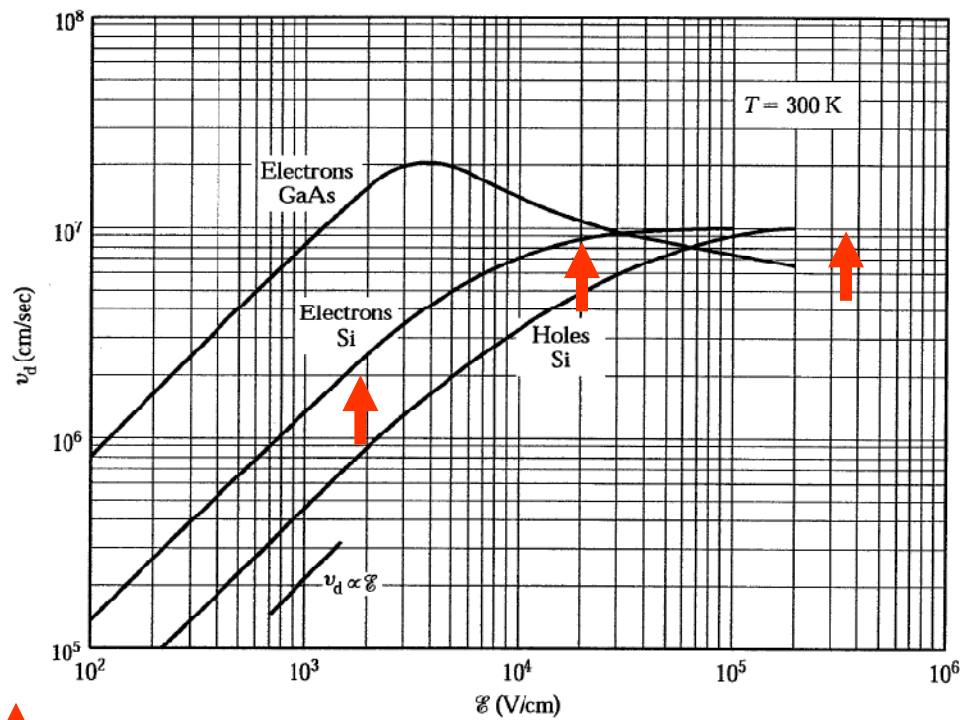
$$v = \frac{q\tau_N}{m_N^*} \mathcal{E}$$



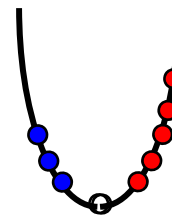
What causes velocity saturation at high fields?

Where does all the mobility formula in device simulator come from?

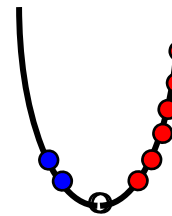
# Velocity Saturation in Si/Ge



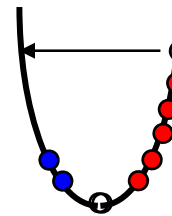
$$\mathcal{E} = 0 \quad J_1 = J^+ - J^- = 0$$



$$\mathcal{E} \ll \mathcal{E}_c \quad J_2 = J^+ - J^- > J_1$$

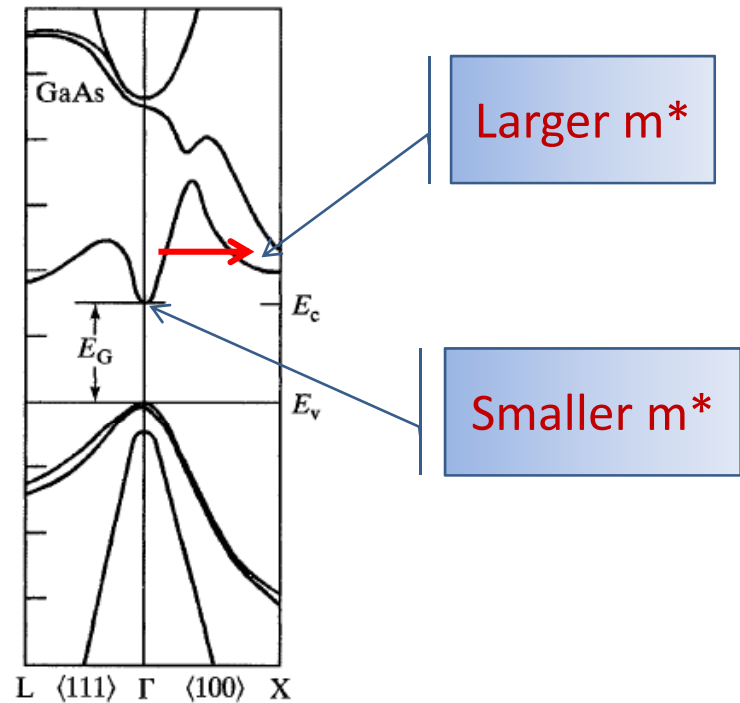
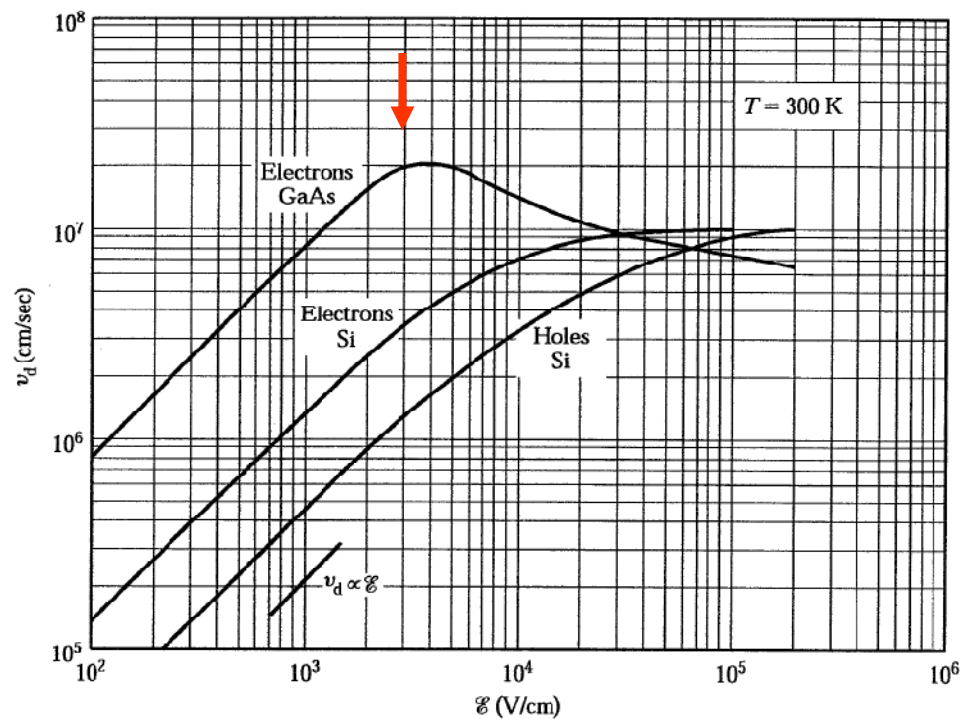


$$\mathcal{E} \approx \mathcal{E}_c \quad J_3 = J^+ - J^- > J_2$$



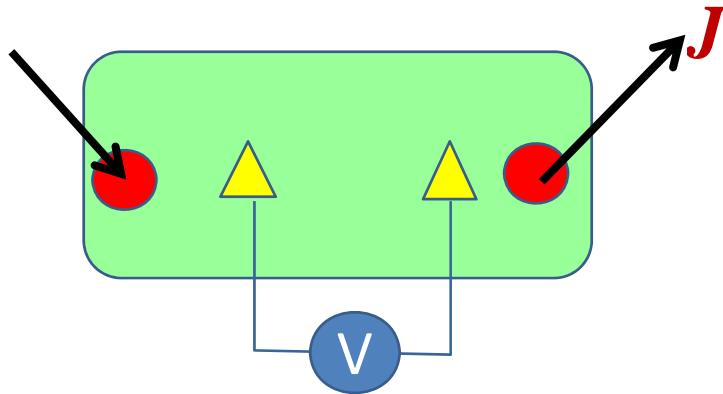
$$\mathcal{E} \gg \mathcal{E}_c \quad J_4 = J^+ - J^- \approx J_3$$

# Velocity Overshoot & Inter-valley Transfer



What type of scattering would you need for inter-valley transfer?

# Doping dependent Resistivity



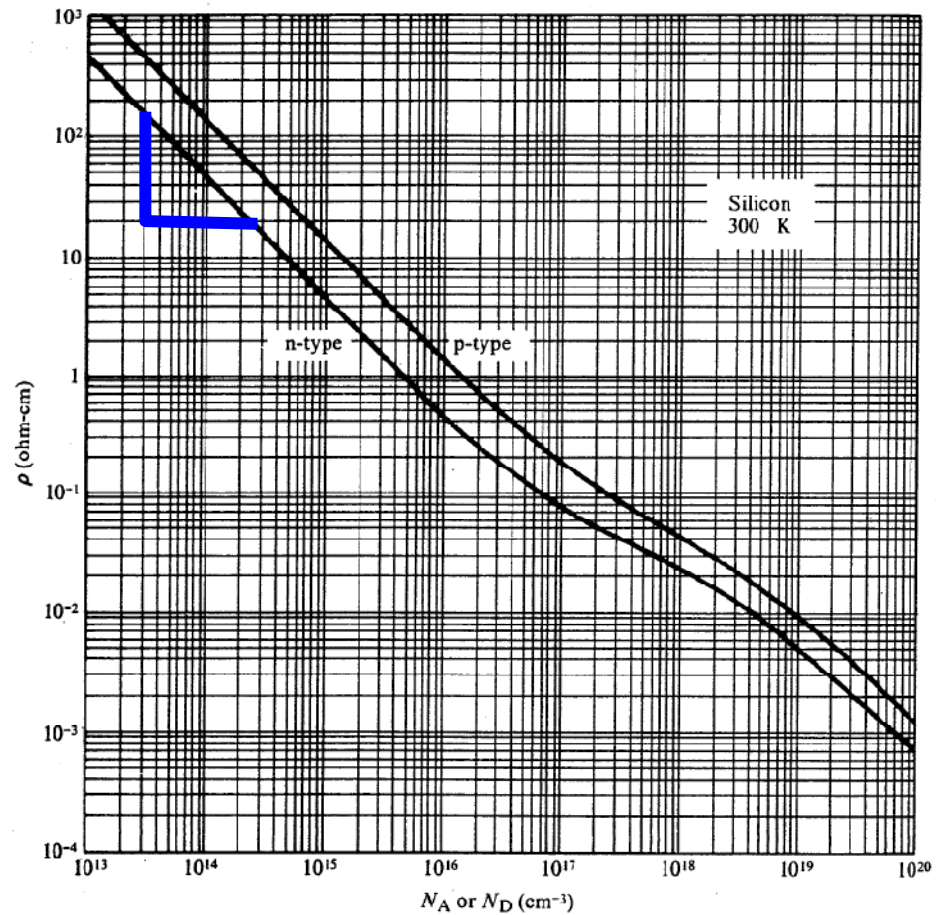
$$\mathcal{E} = \rho J$$

$$J = q(\mu_n n + \mu_p p)\mathcal{E}$$

$$\rho = \frac{1}{q(\mu_n n + \mu_p p)}$$

$$= \frac{1}{q\mu_n N_D} \dots \text{for n-type}$$

$$= \frac{1}{q\mu_p N_A} \dots \text{for p-type}$$



## Conclusion

- 1) Poisson and drift-diffusion equations form a complete semi-classical transport model that can explain wide variety of device phenomena.
- 2) Drift current results from response of electrons/holes to electric field. The physics of mobility is complex and material dependent.
- 3) Constancy of low-field mobility can be checked by experiments.