## NCN

## EE-606: Solid State Devices Lecture 2: Geometry of Periodic Crystals

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## Outline

1) Volume \& surface issues for BCC, FCC, Cubic lattices
2) Important material systems
3) Miller indices
4) Conclusions

Reference: Vol. 6, Ch. 1 (pages 10-17)
Helpful software: Crystal Viewer in ABACUS tool at nanohub.org

## The story so far ...

1. We want to compute current in a material, $\mathrm{J}=\mathrm{q} \mathrm{n} \vee$
2. Computing ( n ) and ( v ) requires knowledge of composition and periodicity of the material
3. If the atoms are arranged in a periodic array, then calculation of ( $n$ ) and (v) will be easier.

Therefore, we are studying the symmetry of the unit cell, hoping that we will be able to identify them in various materials.

Important Bravais lattices
Polonium84

|  | Triclinic | Monoclini | Orthorhomb | Tetragonal | Rhombohedral | Hexagonal | Cubic |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| P | $\overbrace{\alpha}^{\alpha, \beta, \gamma \neq 90^{\circ}}$ |  |  |  |  |  | $\square$ |
| 1 |  |  |  |  |  |  | $\stackrel{\square}{a}$ |
| F |  |  |  |  |  |  | $\underbrace{2}$ |
| C |  |  |  |  |  |  |  |

## Cubic Lattice: Volume Issues



Points per cell
=1/8 points/corner x 8 corners
=1 Point/cell.
(depends on definition of cell)

Number density<br>$=\left(1 / a^{3}\right)$ points $/ \mathrm{cm}^{3}$.<br>(does not depend on cell definition)

## Cubic Lattice: Volume Issues



> Packing density
> =volume filled/total volume
> $R=a / 2$
> $P=(1 / 8) \times(4 / 3) \pi R^{3} \times(8$ corners $) / a^{3}$
> $=\pi / 6$
(does not depend on cell definition)

## Cubic Lattice: Surface Issues



## BCC and FCC lattices



## Hexagonal Closed-Packed

Points per cell


1/2 x 2 @faces =1


1/2×1/3 x 12 @corners =2
3 points/cell

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## Geometry of Lattice Points

In a Bravais lattice, every point has the same environment as every other point (same number of neighbors, next neighbors, ...)
b

$$
R=h \vec{a}+k \vec{b}
$$



Non-Bravais lattice

Bravais lattice with a basis


## Rock-Salt as FCC lattice

$$
=0
$$



For more discussion, see Kittle and Ashcroft/Mermin

## Zinc-Blende FCC Lattice for GaAs



Atoms/cell=(1/8) x8 + (1/2) x6 $+4 \neq 8$
FCC Lattice with a basis
Tetrahedral structure


## Diamond FCC Lattice for Silicon



## Surface Reconstruction



## Hexagonal Closed-Packed for CdS

Focus on (Cd) ...

(Cd) atoms/cell= $(1 / 6) \times 12+(1 / 2) \times 2+3=6$

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## Miller-Indices and Definition of Planes



## Miller Indices: Rules



1. Set up axes along the edges of unit cell
2. Normalize intercepts .... 2, 3, 1
3. Invert/rationalize intercepts ... $1 / 2,1 / 3,1$

$$
3 / 6,2 / 6,6 / 6
$$

4. Enclose the numbers in curvilinear brackets (326)

## Few more rules ...

Negative Intercept


2, 3, -2
$1 / 2,1 / 3,-1 / 2$
$3, \quad 2,-3$
(3 $2 \overline{3}$ )

Intercept at infinity

$\left.\begin{array}{|lll|}\hline 2, & 3, & \propto \\ 1 / 2, & 1 / 3, & 0 \\ 3, & 2, & 0 \\ (3 & 2 & 0\end{array}\right)$

## Bravais-Miller Indices



First three indices sum to zero.

## Where does Miller Indices come from ?

Miller indices: (326)


Normal to the surface
$R_{2} \times R_{1}=\left|\begin{array}{ccc}a & b & c \\ -2 & 3 & 0 \\ -2 & 0 & 1\end{array}\right|=3 a+2 b+6 c$

Vector indices same as Miller indices !
(326) vs. [326]

## Angle between Two Planes



Unit vector normal to plane 1:

$$
N_{1}=\left(h_{1} \vec{a}+k_{1} \vec{b}+l_{1} \vec{c}\right) /\left(h_{1}^{2}+k_{1}^{2}+l_{1}^{2}\right)^{1 / 2}
$$

Unit vector normal to plane 2:

$$
N_{2}=\left(h_{2} \vec{a}+k_{2} \vec{b}+l_{2} \vec{c}\right) /\left(h_{2}^{2}+k_{2}^{2}+l_{2}^{2}\right)^{1 / 2}
$$

$$
\begin{aligned}
& \operatorname{Cos}(\theta)=N_{1} \bullet N_{2} \\
& =\left(h_{2} h_{1}+k_{2} k_{1}+l_{2} l_{1}\right) /\left(h_{2}^{2}+k_{2}^{2}+l_{2}^{2}\right)^{1 / 2}\left(h_{2}^{2}+k_{2}^{2}+l_{2}^{2}\right)^{1 / 2}
\end{aligned}
$$

## Examples ...



$$
\begin{aligned}
& \cos (\theta)=(1 \times 0+0 \times 1+0 \times 1) /(\sqrt{ } 1 \times \sqrt{ } 2)=0 \\
& \text { so } \theta=90 \text { degrees } \\
& (011) \text { surface is normal to }(100) \text { surface }
\end{aligned}
$$

## Example use of Miller Indices .. Find [021] Direction


[100]

(100) $n$-type
$\downarrow \quad \downarrow \quad \downarrow \quad \downarrow \quad \downarrow \quad \downarrow$
$\cos \left(\theta_{1}\right)=(1 \times 0+0 \times 2+0 \times 1) /(1 \times \sqrt{ } 5)=0$, so $\theta_{1}=90$ degrees
[021] vector lies on (100) plane. $\cos \left(\theta_{2}\right)=(0 \times 0+2 \times 1+1 \times 1) /(\sqrt{ } 5 \sqrt{ } 2)$, so $\theta_{2}=18.5$ degrees with respect to [011] direction.

## Conclusions

1. To understand transport in semiconductors, we need to know carrier density ( n ) and carrier velocity ( v ). In order to find these quantities, we need to understand the chemical composition and atomic arrangements.
2. Crystalline material can be built by repeating the basic building blocks. This simplifies the quantum solution of the material, which will allow us to compute n and v for these systems easily.
3. Silicon, GaAs, PbS do not have simple Bravais lattice; but they have Bravais lattice with basis.
4. Often we need to calculate the direction of crystal planes because material properties differ along different planes. Miller indices are one useful way of characterizing crystal planes. In is useful to to review some identities of vector calculus to such calculations involving crystal planes.

## Looking ahead

Number of atoms/volume from crystal structure

Number of electrons available for conduction

$$
n \neq \rho \times N
$$

Number of electrons/atoms

All electrons may be created equally, but they appear do not behave identically!

