

Engineering Space for Light with Metamaterials

Part 1: Electrical and Magnetic Metamaterials

**Part 2: Negative-Index Metamaterials, NLO,
and super/hyper-lens**

Part 3: Cloaking and Transformation Optics

Outline

- What are metamaterials?
- Early electrical metamaterials
- Magnetic metamaterials
- **Negative-index metamaterials**
- Chiral metamaterials
- Nonlinear optics with metamaterials
- Super-resolution
- Optical cloaking

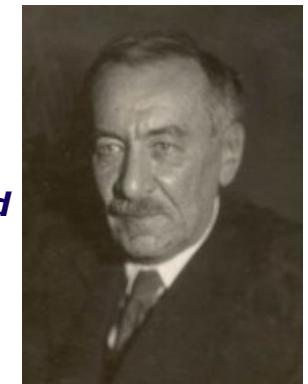
Negative refractive index: A historical review



Sir Arthur Schuster

Sir Horace Lamb

... energy can be carried forward at the group velocity but in a direction that is anti-parallel to the phase velocity...
Schuster, 1904



Negative refraction and backward propagation of waves
Mandel'stam, 1945



V. G. Veselago

Left-handed materials: the electrodynamics of substances with simultaneously negative values of ϵ and μ

Veselago, 1968

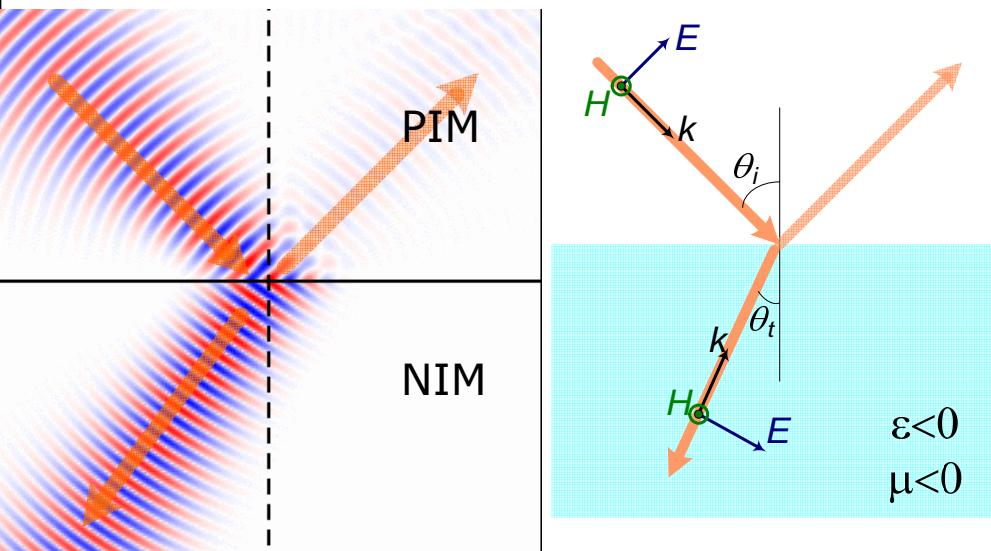


Pendry, the one who whipped up the recent boom of NIM researches

*Perfect lens (2000)
EM cloaking (2006)*

Sir John Pendry

Metamaterials with Negative Refraction



Refraction:

$$n^2 = \epsilon\mu$$

$$n = \pm \sqrt{\epsilon\mu}$$

Warning! Negative refraction ≠

negative refractive index

(e.g. see Peccianti & Assanto: OE, 2007)

$n < 0$, if $\epsilon'|\mu| + \mu'|\epsilon| < 0$

Single-negative:

$n < 0$ when $\epsilon' < 0$ whereas $\mu' > 0$
(F is low)

Double-negative:

$n < 0$ with both $\epsilon' < 0$ and $\mu' < 0$
(F can be large)

Figure of merit

$$F = |n'|/n''$$

Negative Refractive Index in Optics: State of the Art

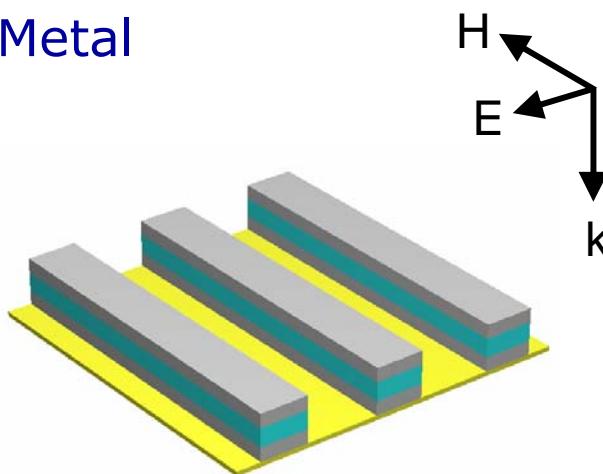
<i>Year and Research group</i>	<i>1st time posted and publication</i>	<i>Refractive index, n'</i>	<i>Wavelength λ</i>	<i>Figure of Merit $F= n' /n''$</i>	<i>Structure used</i>
<u>2005:</u>					
Purdue	April 13 (2005) arXiv:physics/0504091 Opt. Lett. (2005)	-0.3	1.5 μm	0.1	Paired nanorods
UNM & Columbia	April 28 (2005) arXiv:physics/0504208 Phys. Rev. Lett. (2005)	-2	2.0 μm	0.5	Nano-fishnet with round voids
<u>2006:</u>					
UNM & Columbia	J. of OSA B (2006)	-4	1.8 μm	2.0	Nano-fishnet with round voids
Karlsruhe & ISU	OL. (2006) OL (2007)	-1 -1	1.4 μm 1.4 μm	3.0 2.5	Nano-fishnet 3-layer nanofishnet
Karlsruhe & ISU	OL (2006)	-0.6	780 nm	0.5	Nano-fishnet
Purdue	OL (2007)	-0.9 -1.1	770 nm 810nm	0.7 1.3	Nano-fishnet

see review: Nature Photonics v. 1, 41 (2007)

Negative permeability and negative permittivity

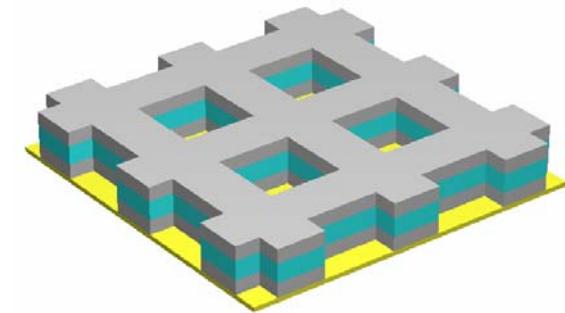
■ Dielectric

□ Metal



Nanostrip pair (TM)
 $\mu < 0$ (resonant)

Nanostrip pair (TE)
 $\varepsilon < 0$ (non-resonant)

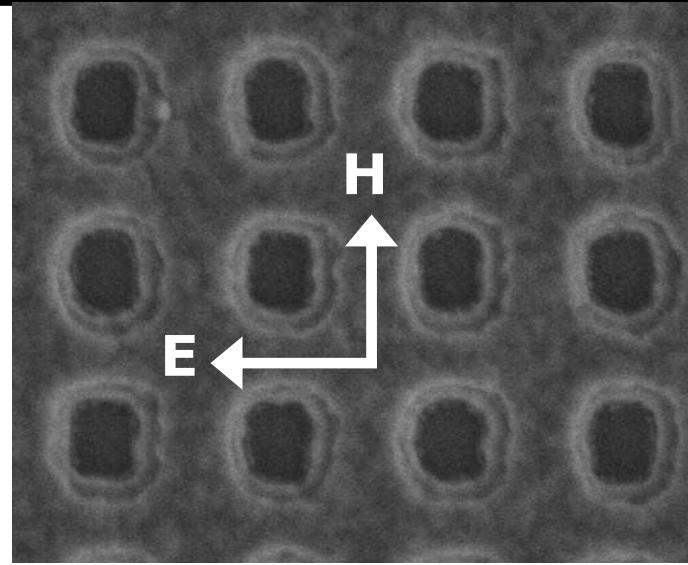


Fishnet
 ε and $\mu < 0$

S. Zhang, et al., PRL (2005)

Sample Geometry (Fishnet Structure)

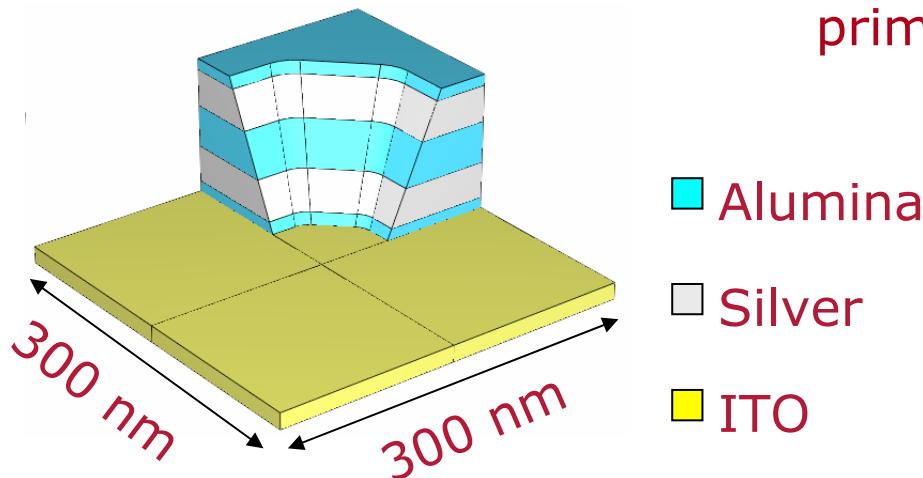
- E-beam lithography
- Period = 300 nm along both axis
- Average width of strips along H = 130 nm
Average width of strips along E = 95 nm



SEM image and
primary polarization

Stacking:

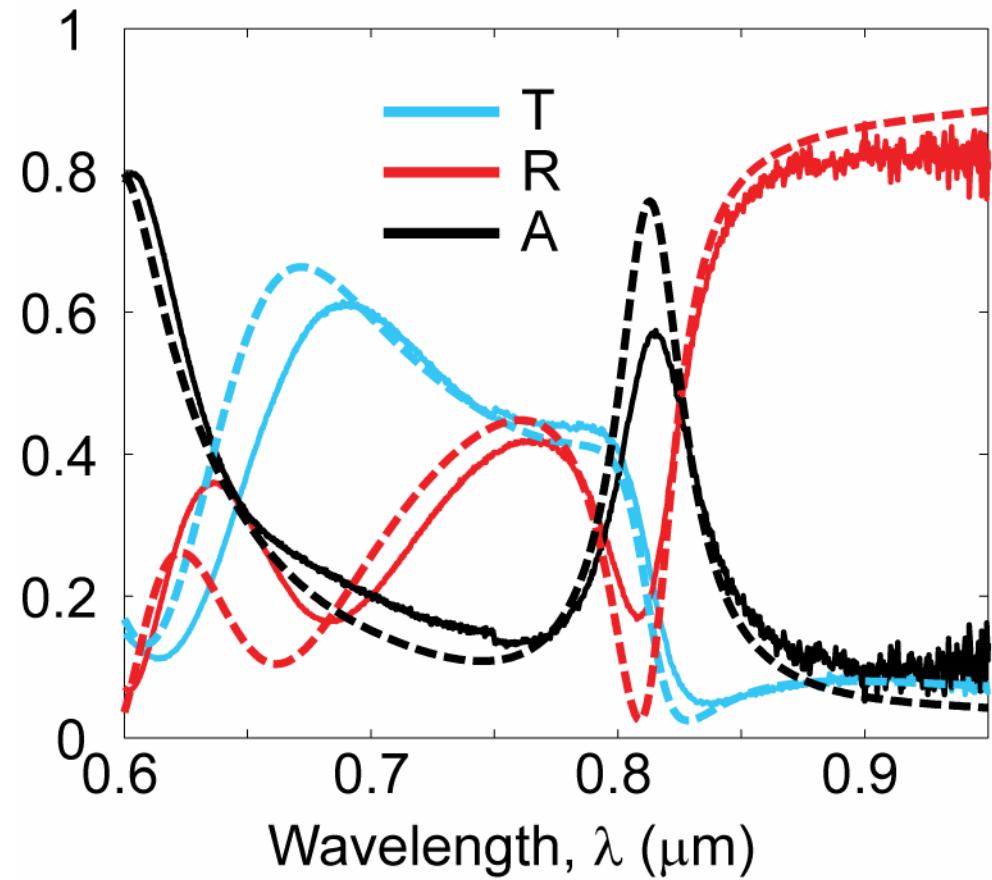
10 nm of Al_2O_3
33 nm of Ag
38 nm of Al_2O_3
33 nm of Ag
10 nm of Al_2O_3



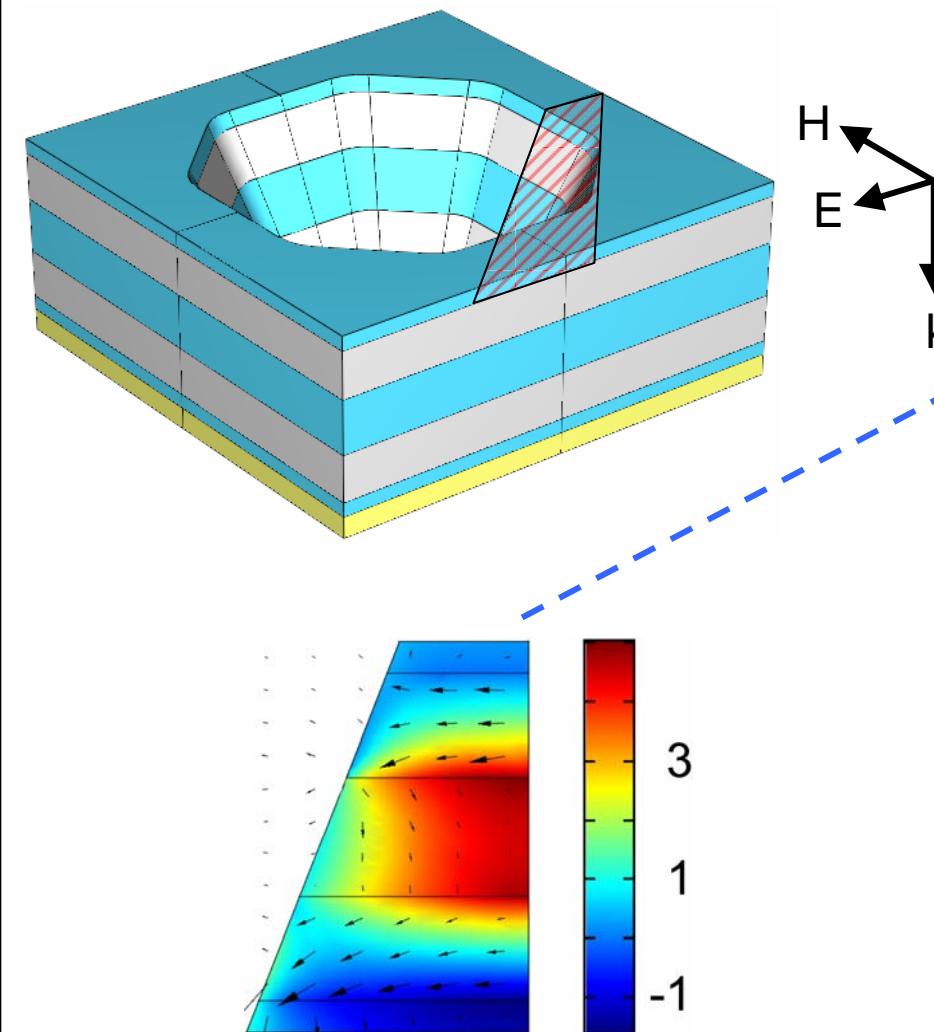
Spectra for Primary Polarization

- Magnetic resonance around $\lambda = 800 \text{ nm}$
- Electric resonance around $\lambda = 600 \text{ nm}$
- Finite Elements

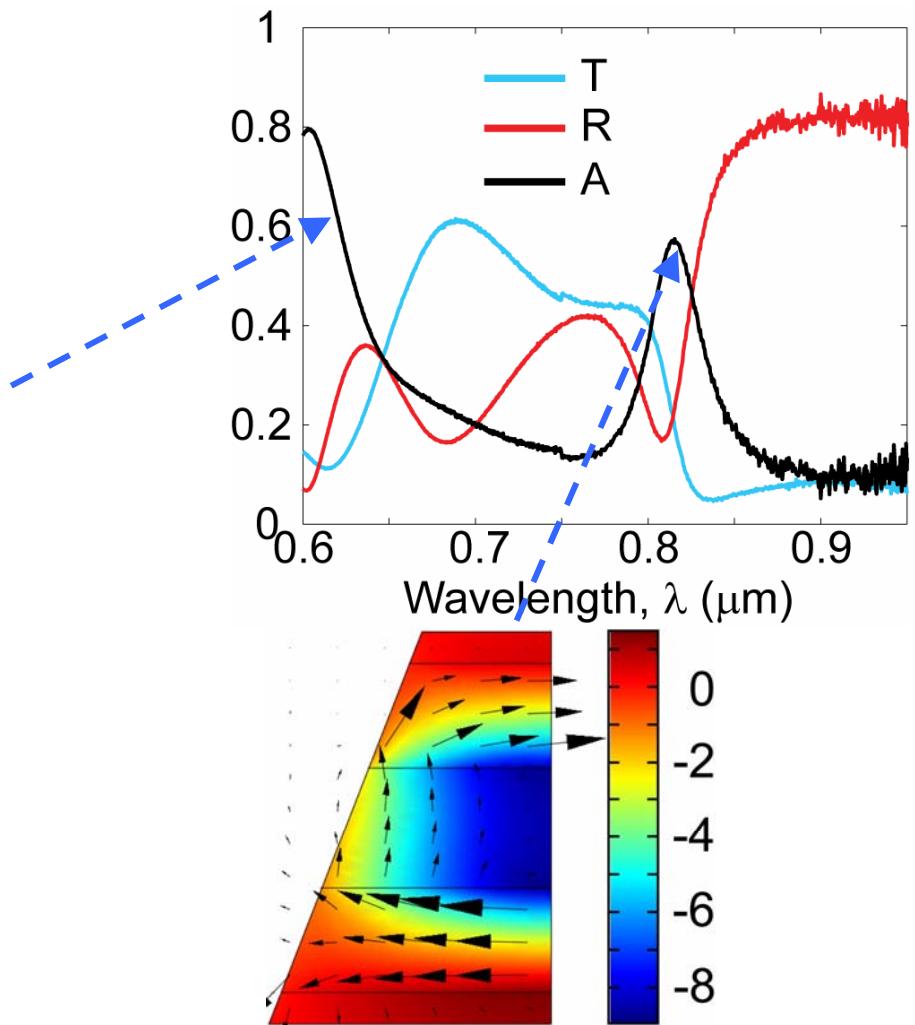
Solid line : Experimental
Dashed line: Simulated



Field Maps for Primary Polarization

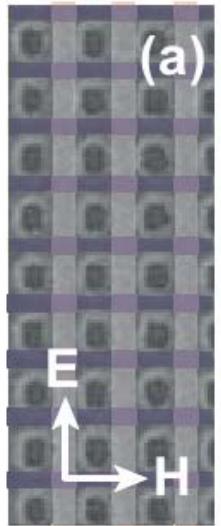
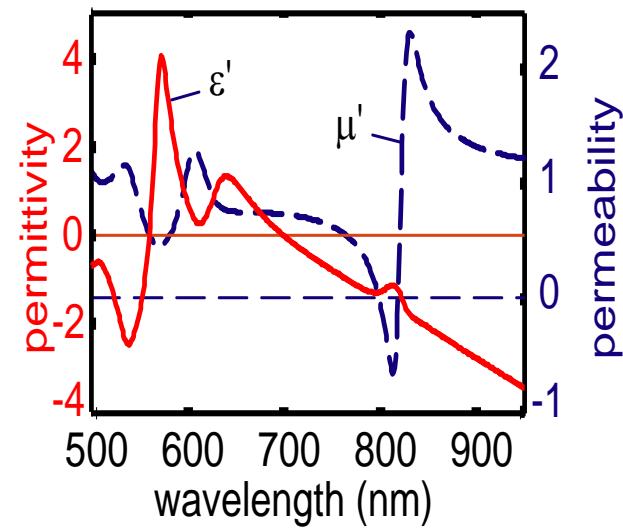
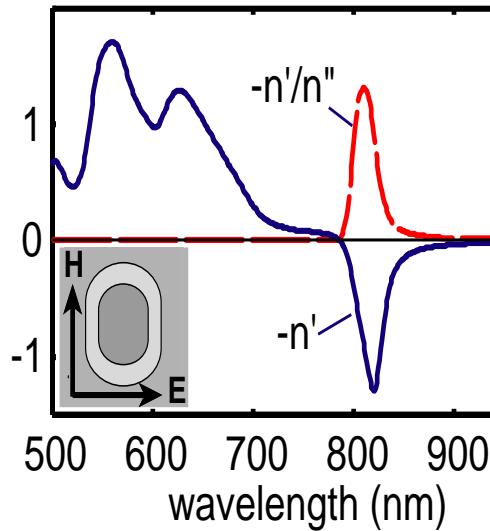
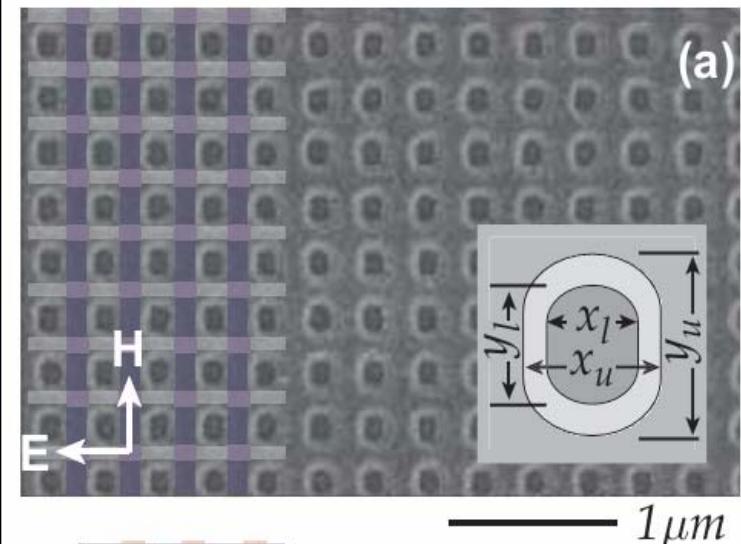


Electrical Resonance, $\lambda = 625 \text{ nm}$

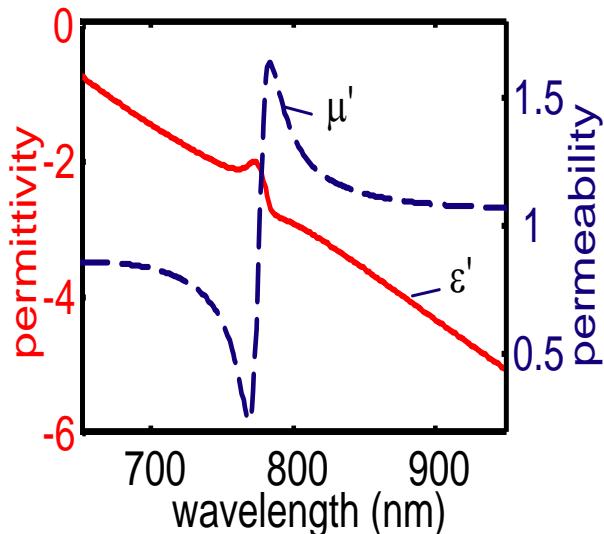
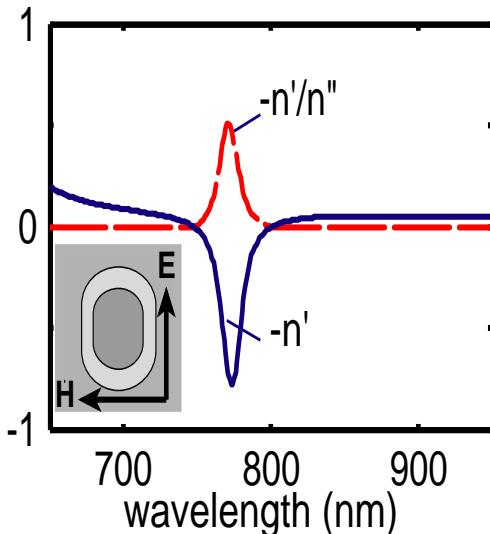


Magnetic Resonance, $\lambda = 815 \text{ nm}$

Double Negative NIM ($n'=-1.0$, FOM=1.3, at 810 nm) Single Negative NIM ($n'=-0.9$, FOM=0.7, at 770 nm)



**Chettiar et al
OL (2007)**

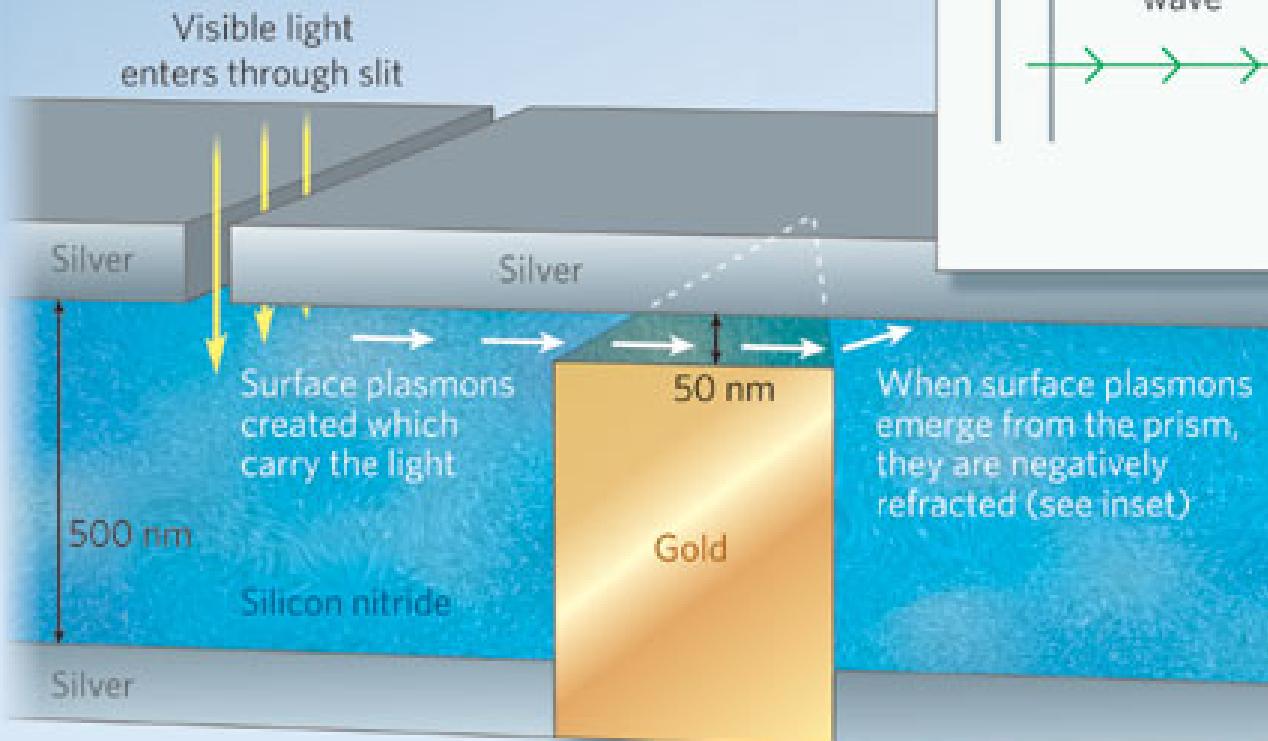


Summary on negative refractive index

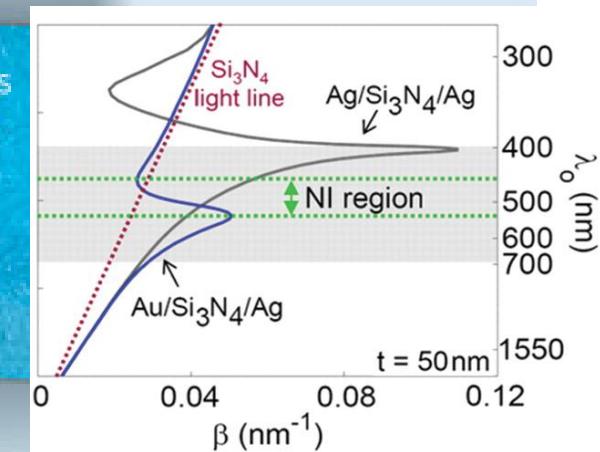
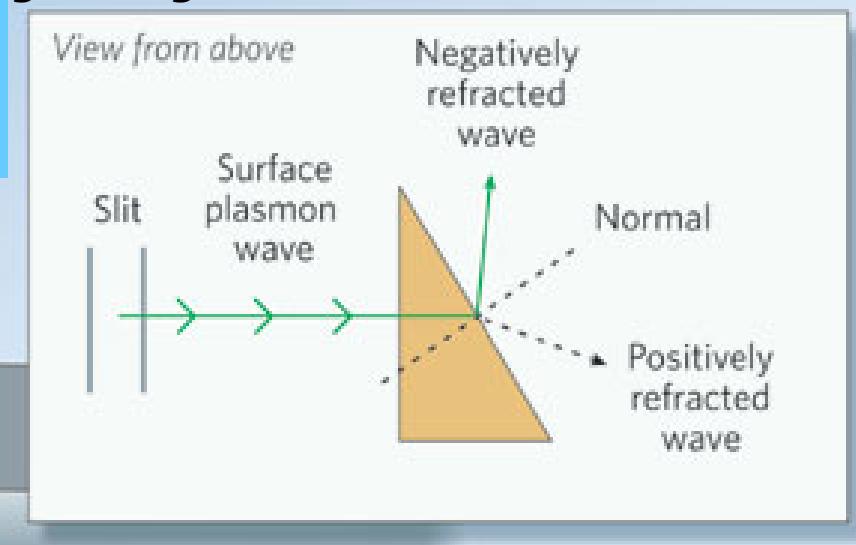
- A Double Negative NIM (Negative index material) is demonstrated at a wavelength of ~ 810 nm
- The sample exhibits a figure of merit ($-n'/n''$) of 1.3 and a transmittance of 25% at 813 nm
- The same sample shows Single Negative NIM behavior for the orthogonal polarization at a wavelength of ~ 770 nm with a figure of merit of 0.7 and a transmittance of 10%

Negative Refraction for Waveguide Modes

negative refraction for 2D SPPs in waveguides



An mode index of ~ -5 is obtained at the green light.

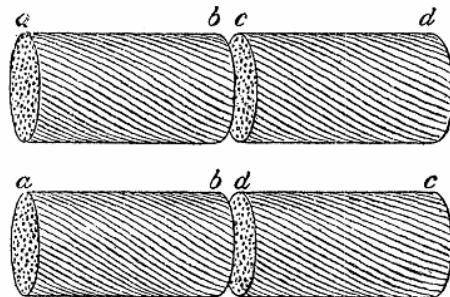


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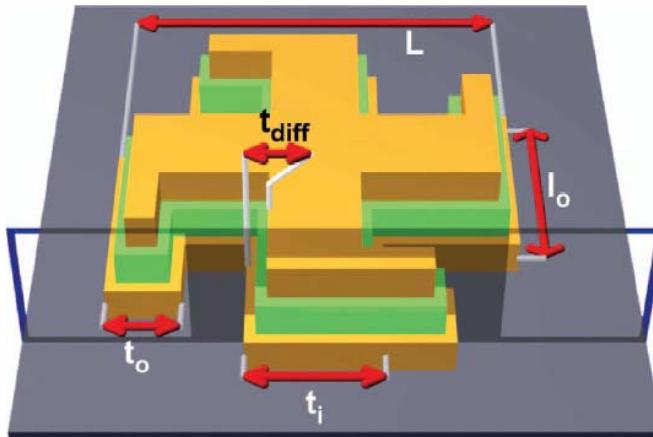
Chiral Optical Elements

Bose's Artificial chiral molecules: Twisted jute elements

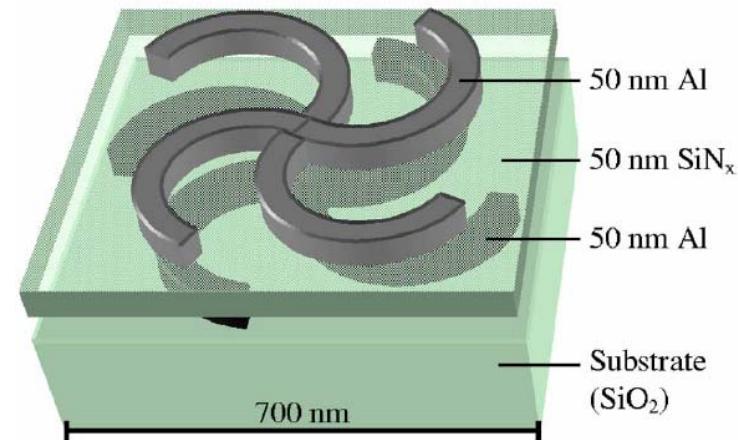


J. C. Bose, *Proceeding of Royal Soc. London*, 1898

Optical counterparts:



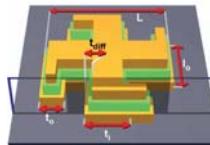
Decher, Klein, Wegener and Linden
Opt. Exp., 2007



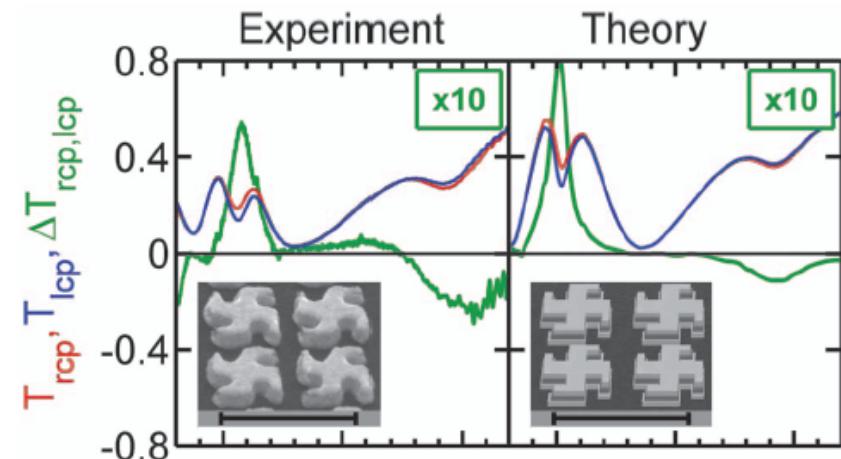
The Zheludev group, U. Southampton
Appl. Phys. Lett., 2007

Chiral Effects in Optical Metamaterials

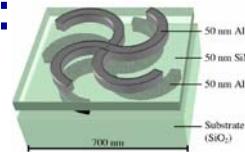
Circular dichroism:



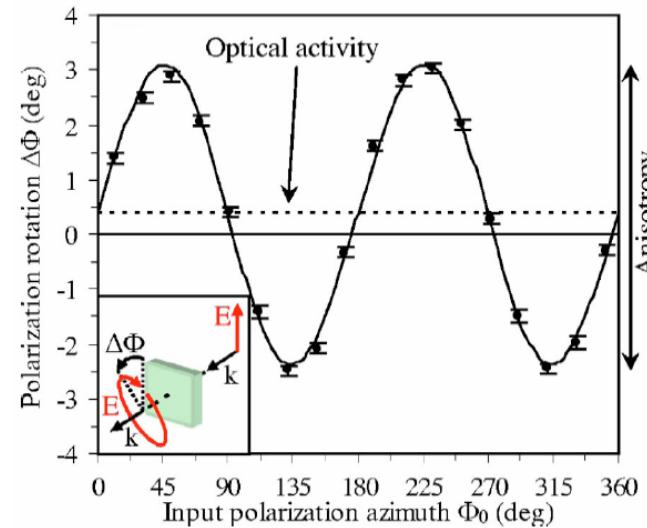
Decher, Klein, Wegener and Linden
Opt. Exp., 2007



Giant optical gyrotropy:



The Zheludev group, U. Southampton
Appl. Phys. Lett., 2007



Chirality can ease obtaining $n < 0$:
Tretyakov, et al (2003), Pendry (Science 2004)

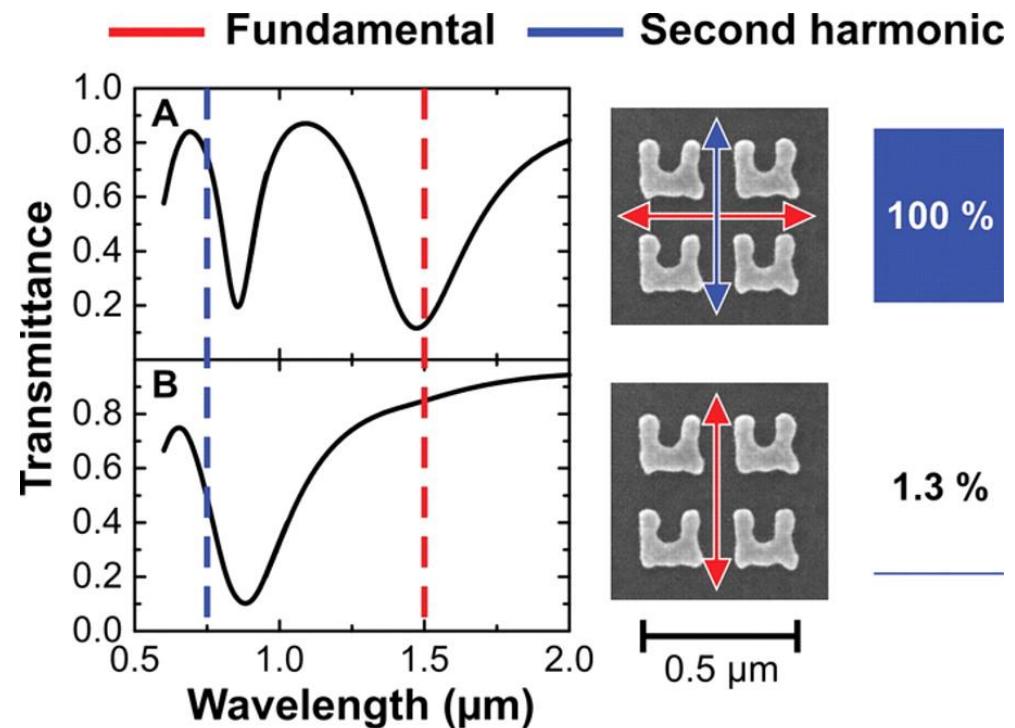
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SHG and THG from Magnetic Metamaterial

Excitation when magnetic resonance is excited (1st pol)

Excitation at 2nd pol.
(no magnetic resonance)



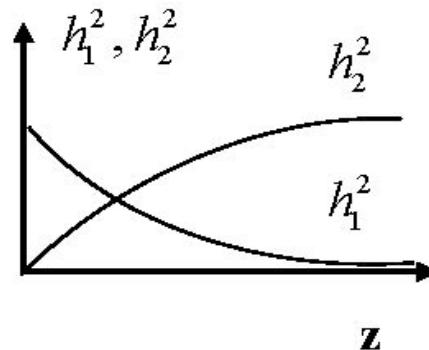
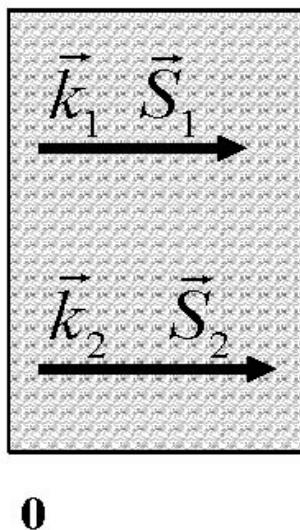
SHG: Klein, Enkrich, Wegener, and Linden, Science, 2006

SHG & THG: Klein, Wegener, Feth and Linden, Opt. Express, 2007

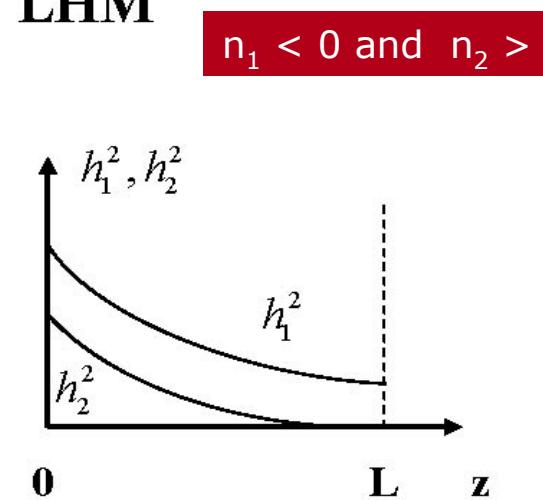
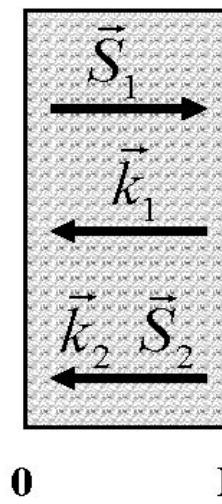
NLO in NIMs: SHG

**Backward Waves in NIMs:
Distributed feedback, cavity-like amplification, etc.**

RHM



LHM



$$n_1 < 0 \text{ and } n_2 > 0$$

Manley-Rowe Relations

$$\frac{dS_1}{dz} - \frac{dS_2}{dz} = 0,$$

Phase-matching: $\epsilon_1 = -\epsilon_2, k_2 = 2k_1 \rightarrow h_1^2(z) - h_2^2(z) = C$

SHG in NIMs: Nonlinear 100% Mirror

100% reflective SHG Mirror !

$$h_1^2(z) - h_2^2(z) = C$$



Finite Slab:

$$\sqrt{C} \kappa L = \arccos(\sqrt{C} / h_{10})$$

$$h_2(z) = \sqrt{C} \tan[\sqrt{C} \kappa(L-z)]$$

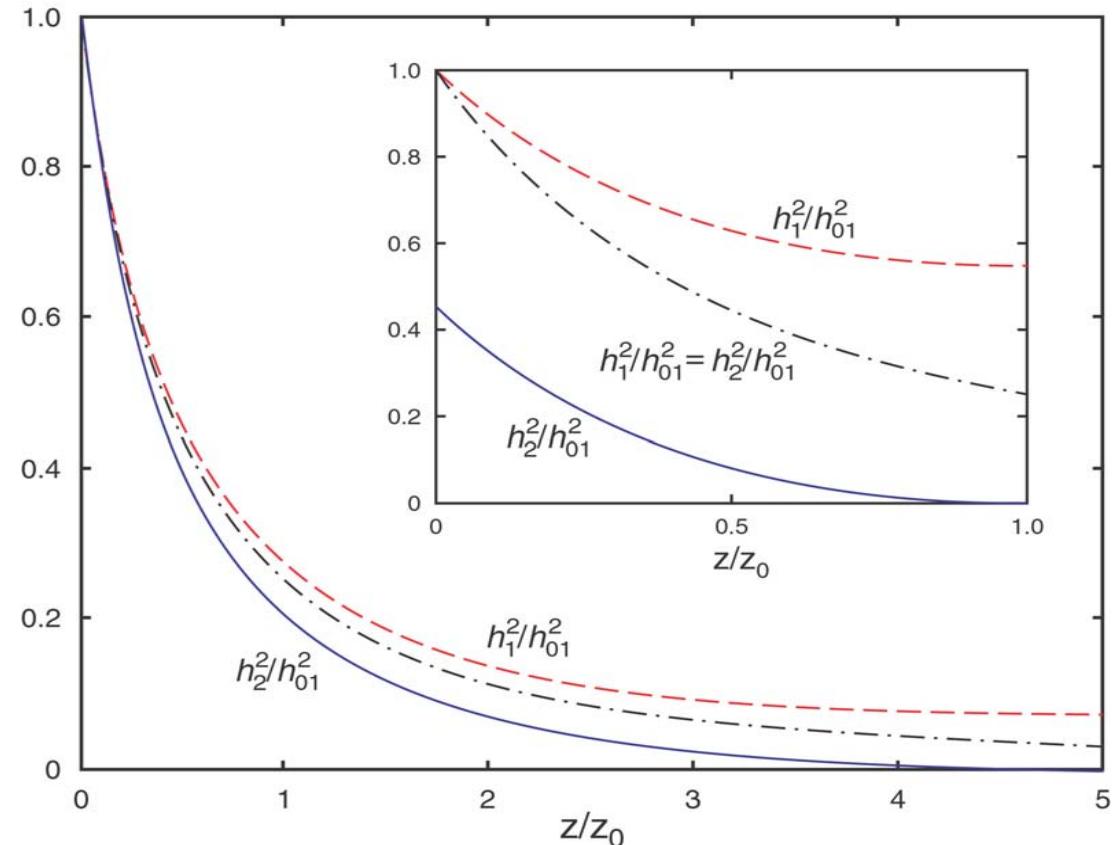
$$h_1(z) = \sqrt{C} / \cos[\sqrt{C} \kappa(L-z)]$$

Semi-Infinite Slab:

$$C = 0, h_2(z) = h_1(z)$$

$$h_2(z) = h_{10} / [1 + (z / z_0)]$$

$$\kappa = 4\pi\chi^{(2)}\varepsilon_2\omega_2^2 / k_2 c^2 \quad z_0 = [\kappa h_{10}]^{-1}$$

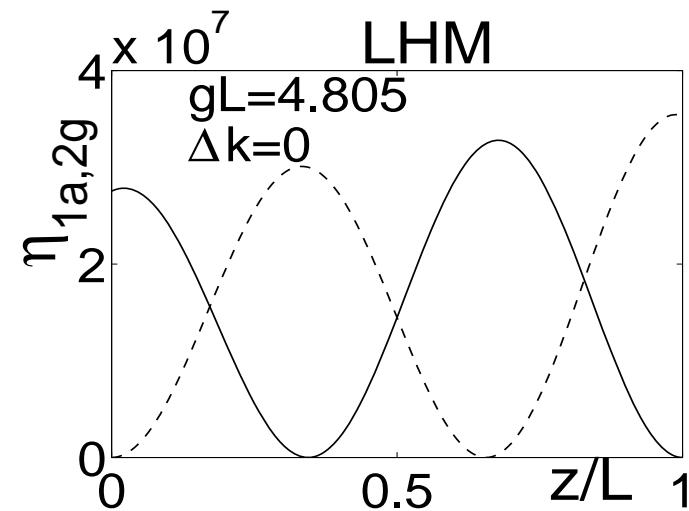
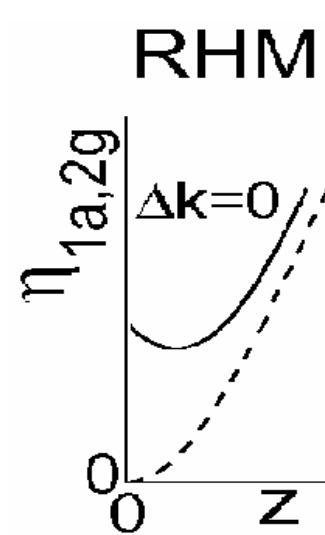
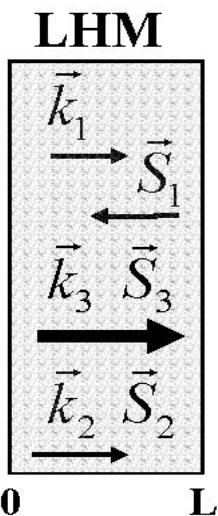
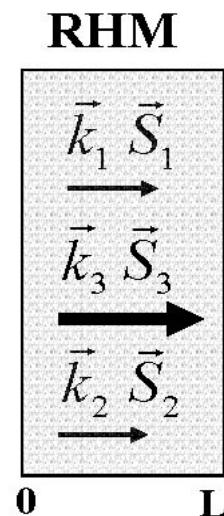


Other work on SHG:
Kivshar et al; Zakhidov et al

Optical Parametric Amplification (OPA) in NIMs

$$\omega_3 = \omega_1 + \omega_2 \quad (n_1 < 0, n_2, n_3 > 0)$$

S_3 - Control Field (pump)



$$\eta_{1a} = |a_1(z)/a_{1L}|^2, \eta_{1g} = |a_1(z)/a_{20}|^2, \eta_{2g} = |a_2(z)/a_{1L}|^2$$

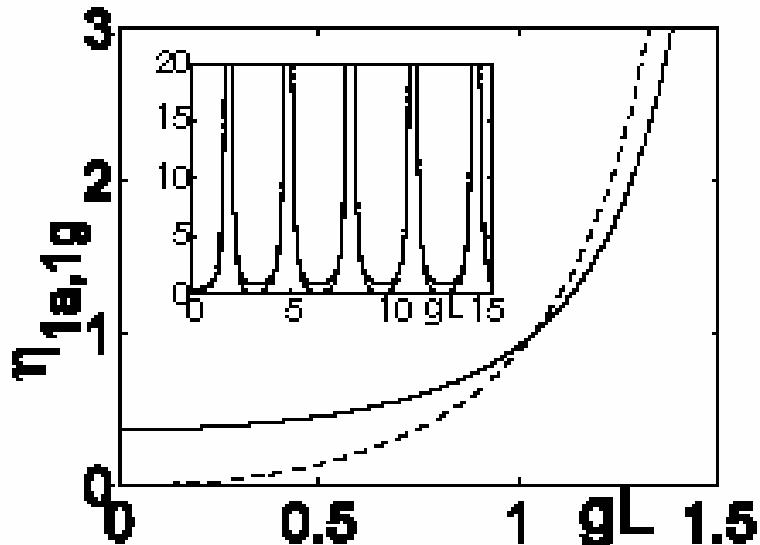
$$g = \left(\sqrt{\omega_1 \omega_2} \sqrt[4]{\epsilon_1 \epsilon_2 / \mu_1 \mu_2} \right) (8\pi/c) \chi^{(2)} h_3$$

Manley-Rowe Relations:

$$\frac{d}{dz} \left(\frac{S_1}{\hbar \omega_1} - \frac{S_2}{\hbar \omega_2} \right) = 0$$

Popov, VMS, Opt. Lett. (2006)
Appl. Phys. B (2006)
For SHG see also Agranovich et al
and Kivshar et al

OPA in NIMs: Loss-Compensator and Cavity-Free Oscillator



$$\Delta k = 0$$

Backward waves in NIMs ->
Distributed feedback & cavity-like
amplification and generation

Popov, VMS, OL (2006)

$$g = \left(\sqrt{\omega_1 \omega_2} \sqrt[4]{\varepsilon_1 \varepsilon_2 / \mu_1 \mu_2} \right) (8\pi/c) \chi^{(2)} h_3 \quad \eta_{1a} = |a_1(z)/a_{1L}|^2, \eta_{1g} = |a_1(z)/a_{20}|^2$$

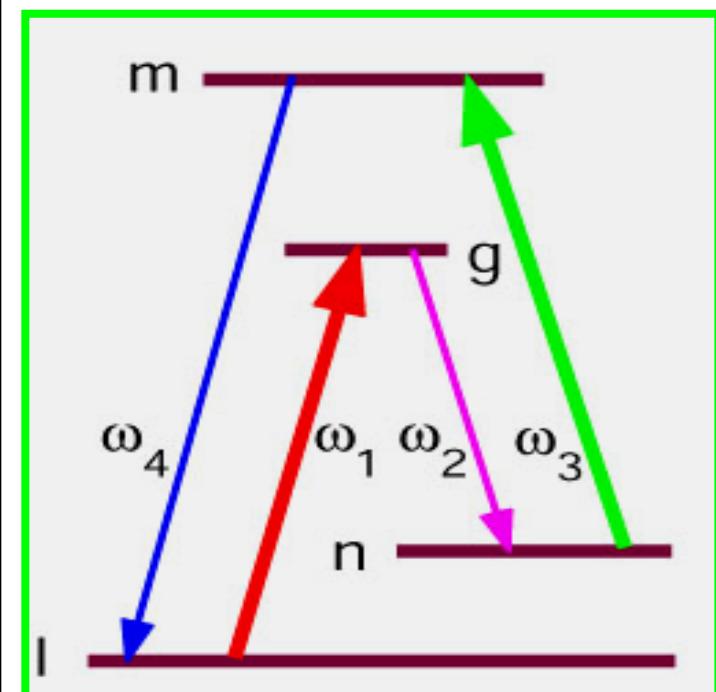
Resonances in output amplification and DFG

$$\alpha_1 L = 1, \alpha_2 L = 1/2$$

- **OPA-Compensated Losses**
- **Cavity-free (no mirrors) Parametric Oscillations**
- **Generation of Entangled Counter-propagating LH and RH photons**

OPA with 4WM

Four-level $\chi^{(3)}$ centers embedded in NIM



$\chi^{(3)}$ -OPA assisted by the Raman Gain:

ω_4 – signal; ω_1, ω_3 – control fields

$$\omega_2 = \omega_1 + \omega_3 - \omega_4 \text{ – idler}$$

(Raman-enhanced; contributes back to OPA at ω_4)

- $\chi^{(3)}$ -OPA: compensation of losses; transparency and amplification at ω_4
- Cavity-free generation of counter-propagating entangled right- and left-handed photons
- Control of local optical parameters through quantum interference

See talk tomorrow by Popov et al on NLO in MMs

Popov, et al OL (2007)

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Super-resolution: Amplification of Evanescent Waves Enables sub- λ Image!

Waves scattered by an object have all the Fourier components

$$k_z = \sqrt{k_0^2 - k_x^2 - k_y^2}$$

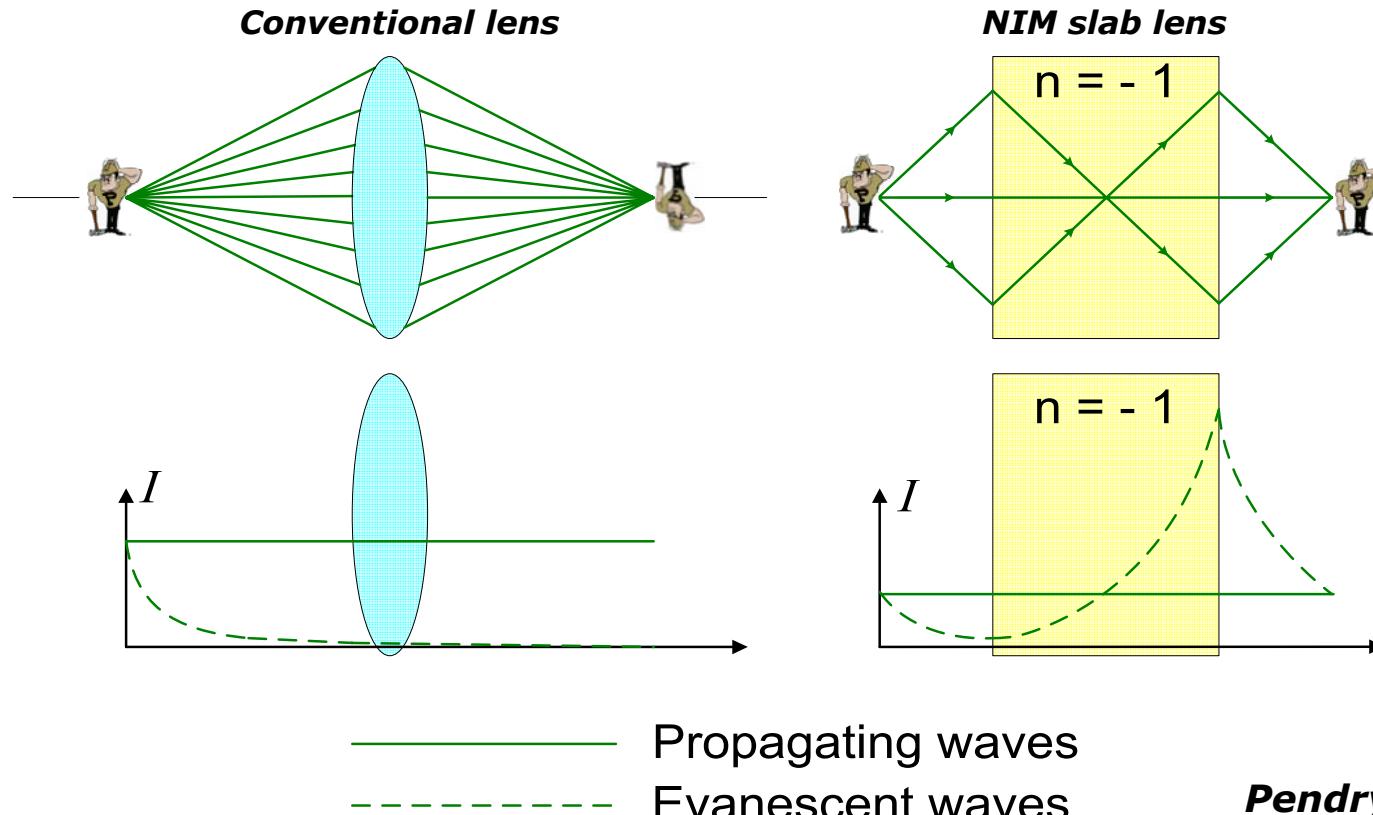
The propagating waves are limited to:

$$k_t = \sqrt{k_x^2 + k_y^2} < k_0$$

To resolve features Δ , we must have

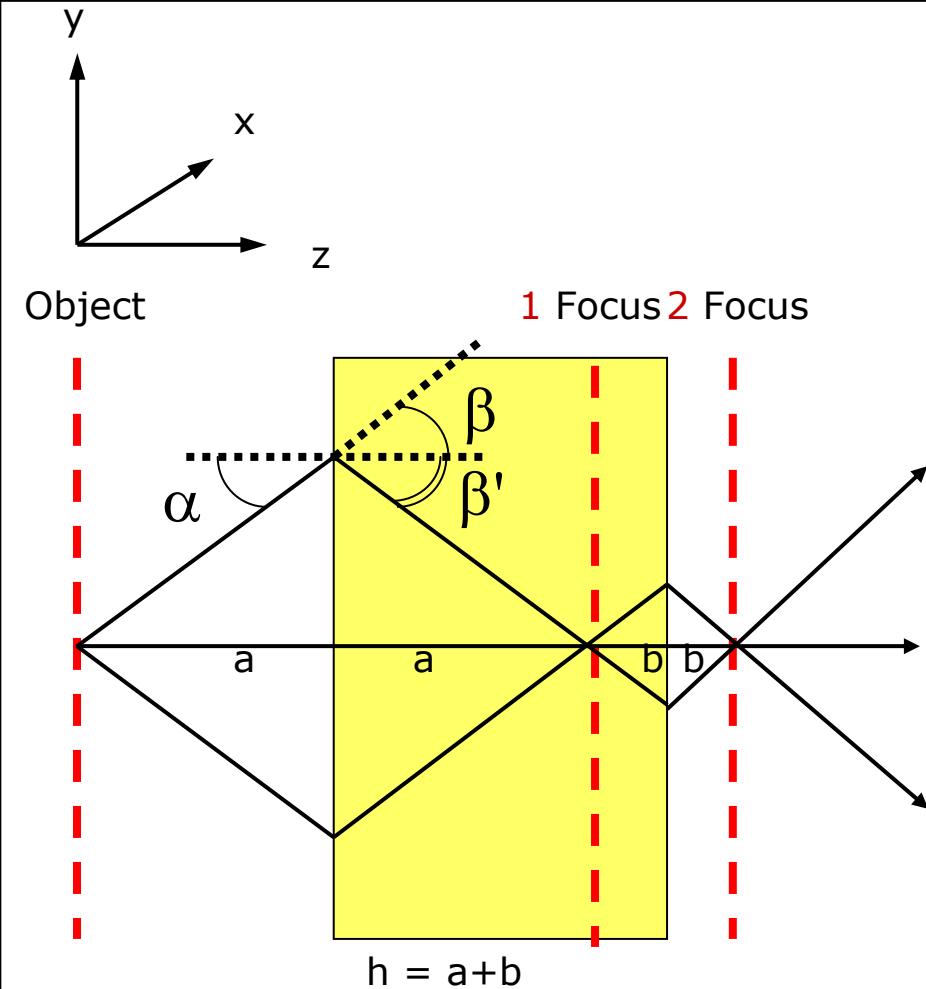
$$\lambda_t = 2\pi/k_t < \Delta, \quad \Delta < \lambda \Rightarrow k_t = \sqrt{k_x^2 + k_y^2} > k_0, \quad k_z^2 < 0$$

The evanescent waves are “re-grown” in a NIM slab and fully recovered at the image plane



Pendry, PRL, 2000

Perfect Lens



$$n = -1$$

$$(\varepsilon = -1; \mu = -1)$$

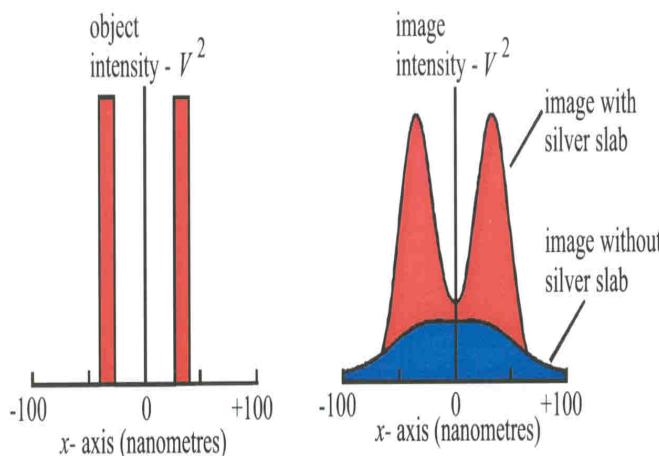
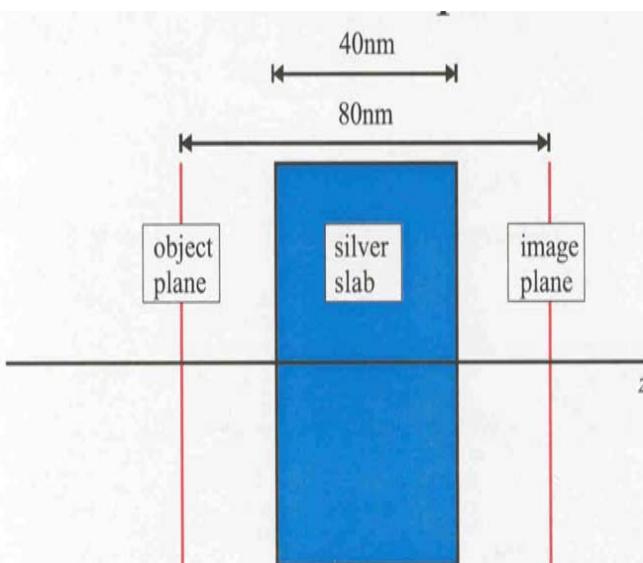
$$\sin(\alpha) = n \sin(\beta') = -\sin(\beta') = \sin(\beta)$$

$$E(y, z) = \sum_q A_q \exp\left(iqy + i\sqrt{(nk)^2 - q^2} z\right)$$

Phase shift =

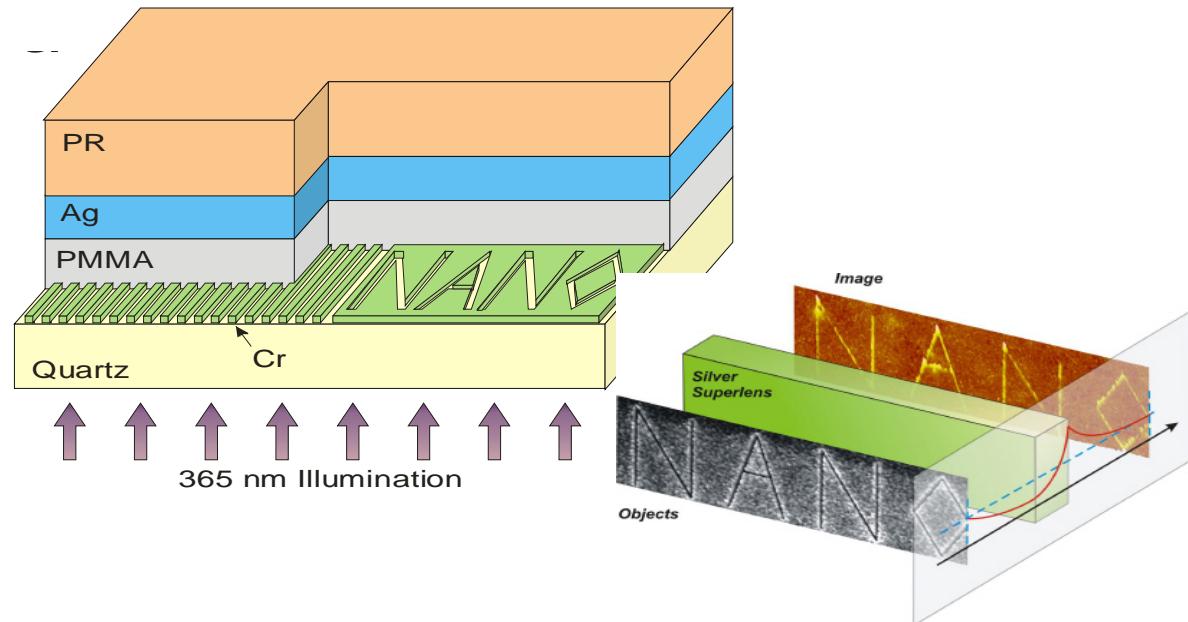
$$\left(iqy + i\sqrt{k^2 - q^2} z\right) + \left(-iqy - i\sqrt{(-k)^2 - q^2} z\right) = 0$$

The Poor Man's (Near-Field) Superlens ($\epsilon < 0$, $\mu = 1$)



Original implementation by Pendry: use a plasmonic material (silver film) to image 10 nm features with $hw = 3.48$ eV;

$$\epsilon = 5.7 - 9^2 / \omega^2 + 0.4i \quad (= -\epsilon_h)$$



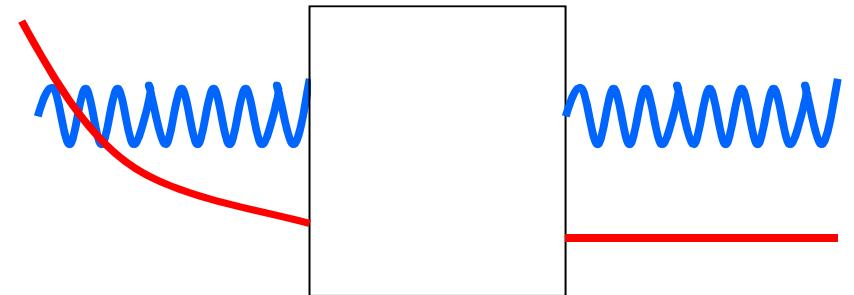
Near-field super-lens (NFSL)

super-resolution with superlens: Zhang et al. (2005); Blaikie, et al (2005)

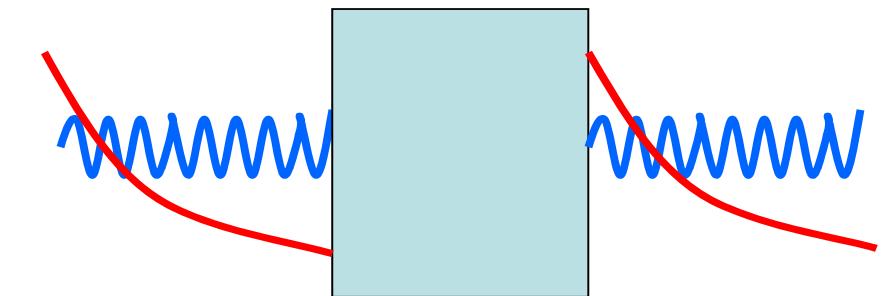
Mid-IR: Shvets et al. (2006)

Superlens High and Low

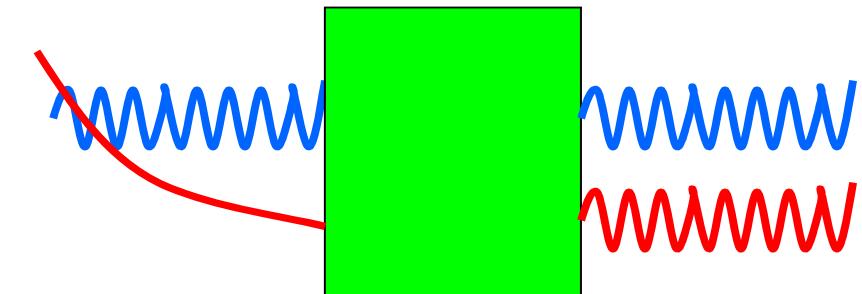
Ordinary Lens:
evanescent field lost



Super Lens:
evanescent field enhanced
but decays away from the lens
* LIMITED TO NEAR FIELD
* EXPONENTIALLY SENSITIVE
TO DISORDER, LOSSES,...

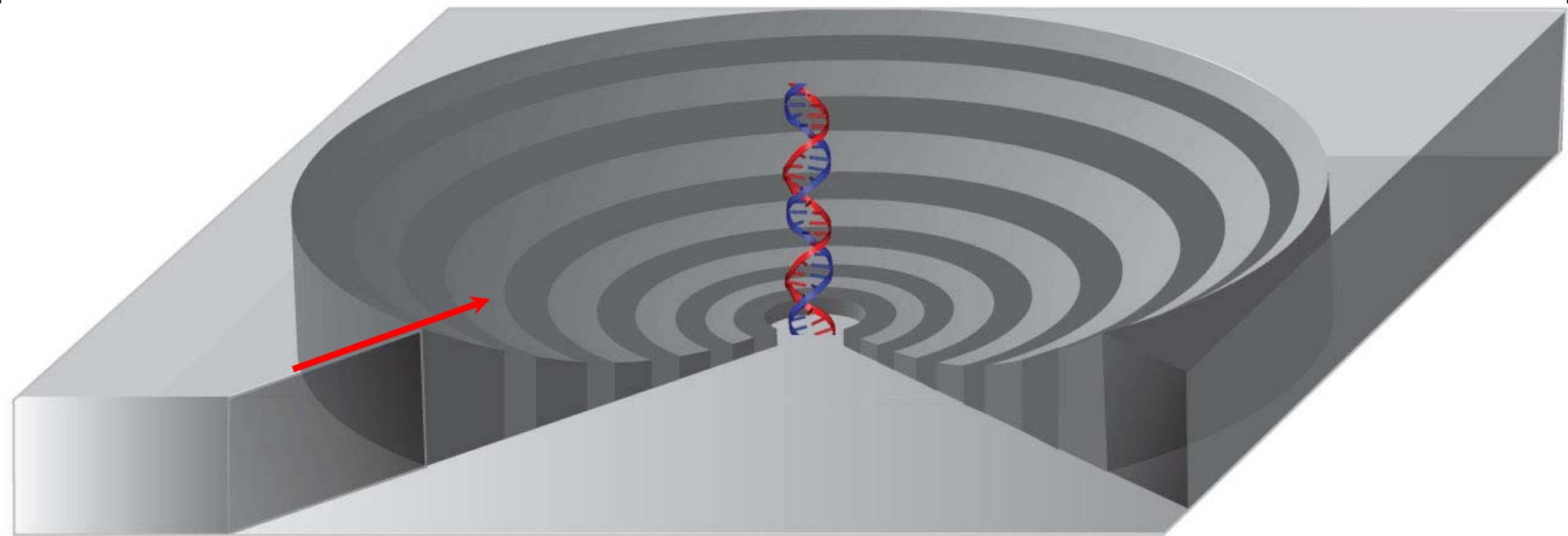


Hyper Lens:
evanescent field converted
to propagating waves (that do
not mix with the others)



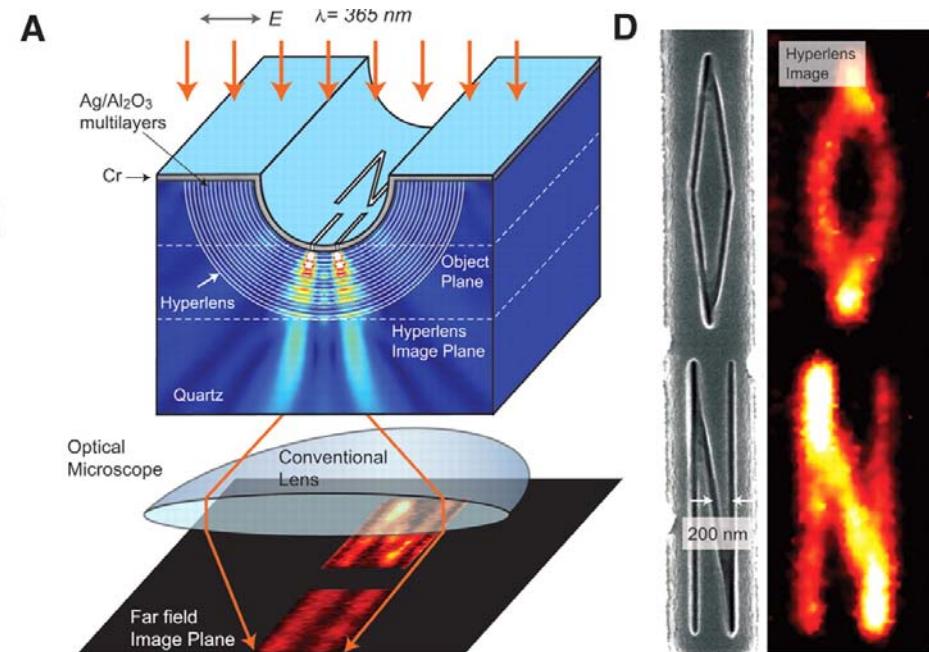
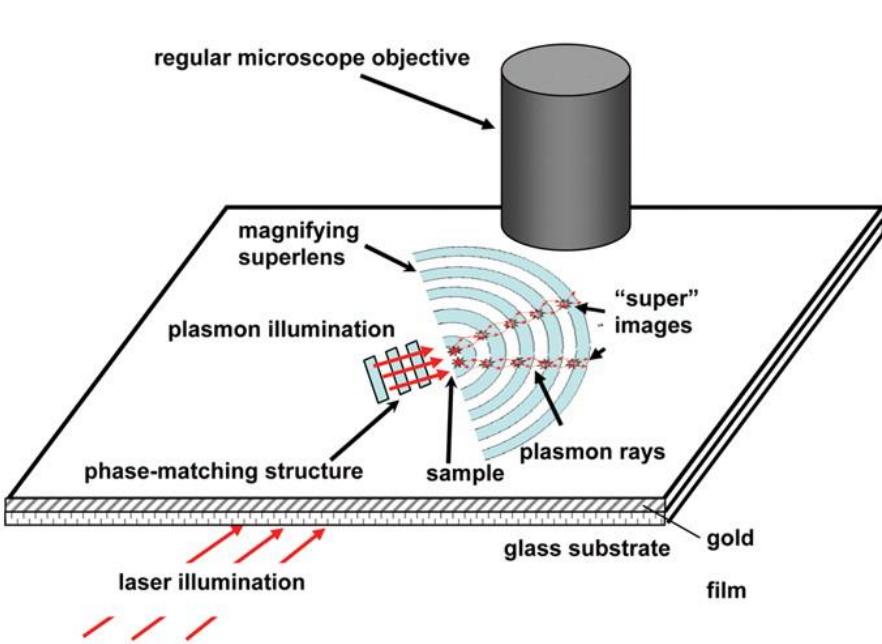
Hyperlens:

Converting evanescent components to propagating waves
(Narimanov eta al; Engheta et al)



Far-field sub- λ imaging

Optical Hyperlens



Theory:

Jacob, Narimanov, OL, 2006

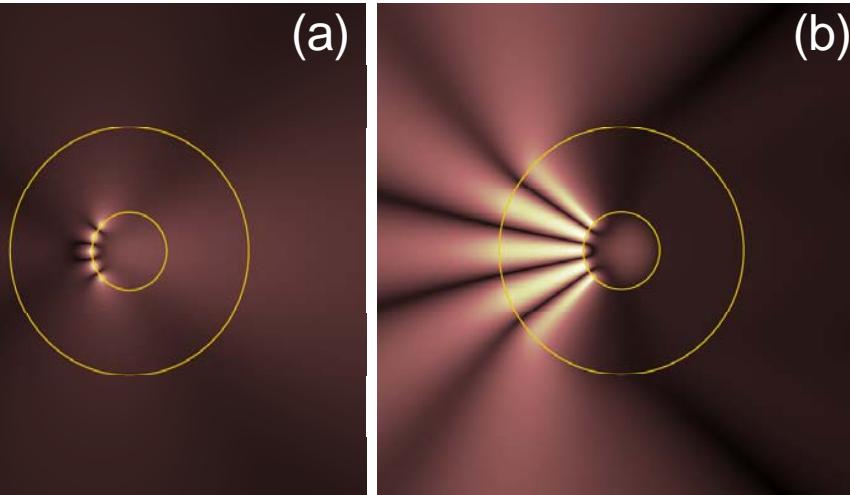
Salandrino, Engehta, PRB, 2006

Experiments:

Z. Liu et al., Science, 2007

Smolyaninov et al., Science, 2007

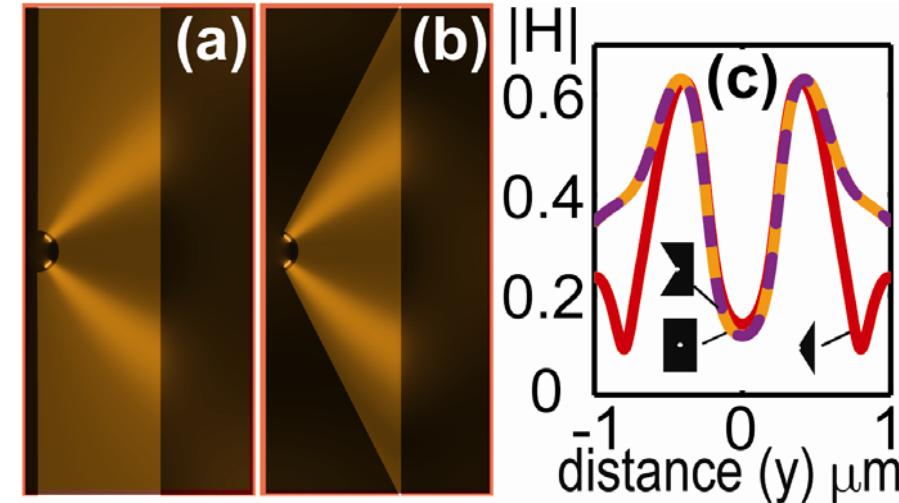
Advanced Optical Hyperlens



Impedance-matched hyperlens

Kildishev, Narimanov

(Opt. Lett., 2007)



Flat hyperlenses:
 $\frac{1}{2}$ - & $\frac{1}{4}$ -body lenses

Kildishev, Shalaev
(Opt. Lett., 2008)

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