Elastic scattering

Lecture 2 Part 2

On to the wave perspective ...

Time independent propagation only one arbitrary direction (+r):

Plane wave:
$$\psi = \psi_o \exp(ikr)$$
 with $k = \sqrt{\frac{2mE}{h^2}}$

Easy to add time dependence back in:

$$\Psi = \Psi_{o} \exp(ikr) \exp\left[-i\frac{E}{h}t\right]$$

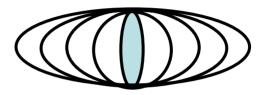
Have a wave, with energy given by:

$$\mathbf{E} = \frac{h^2 \mathbf{k}^2}{2\mathbf{m}}$$
 k $\mathbf{\Theta}$ fixed
p $\mathbf{\Theta}$ infinite (unknown)

Electron in free space

Spherical wave:

 Isotropic wave propagating outwards from a point



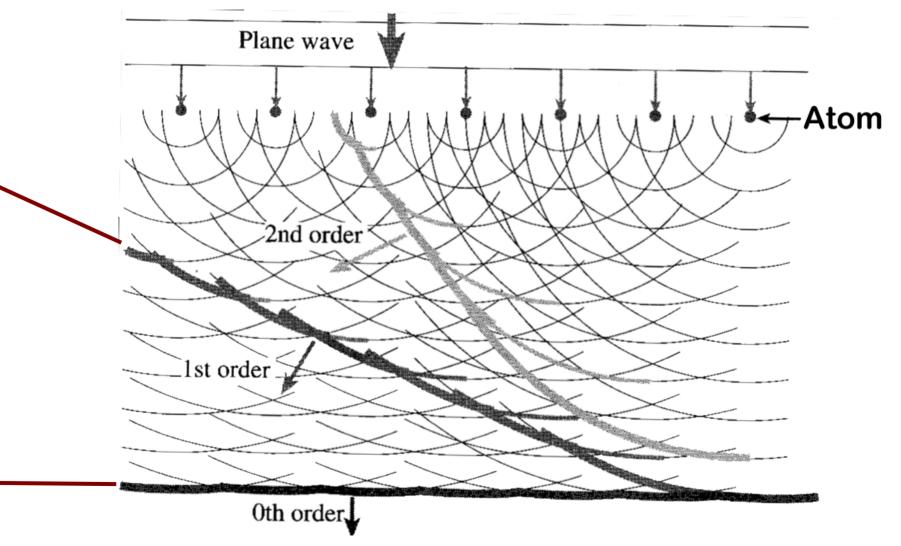
Wave Eqn:

$$\frac{\partial^2 \left(\mathbf{r} \Psi \right)}{\partial t^2} = \mathbf{v}^2 \frac{\partial^2 \left(\mathbf{r} \Psi \right)}{\partial \mathbf{r}^2}$$

Solution: $\Psi = \Psi_{o} \frac{\exp(ikr - \omega t)}{r}$ Time independent

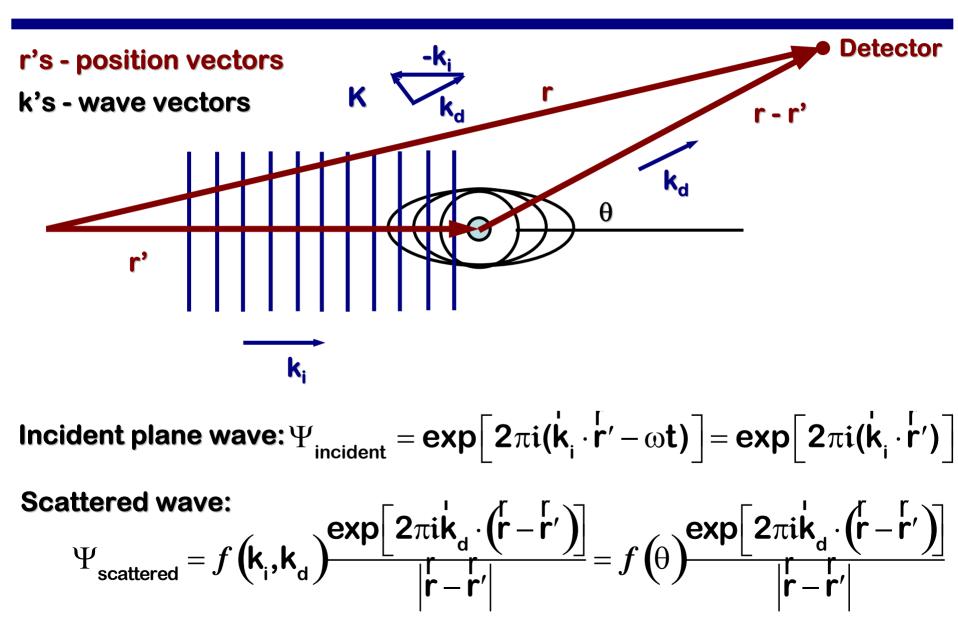
$$\psi = \psi_{o} \frac{\exp(ikr)}{r}$$

Coherence & Incoherence



These waves have the same phase (are coherent) at these angles

Coherent elastic scattering



What is $f(\theta)$?

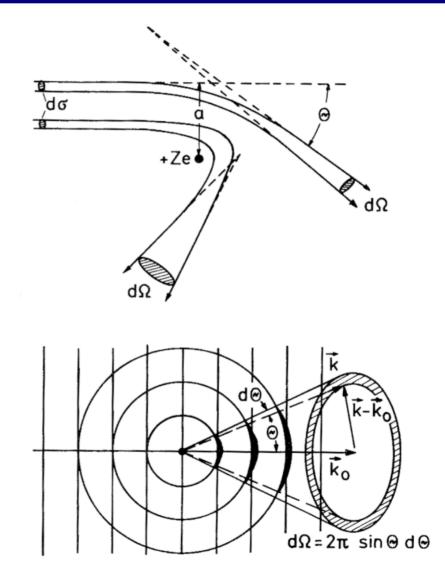
 Measures the strength of the scattering event

Elastic scatter has a strong angular dependence

Remember: "scattering cross section" used to describe strength of scatter earlier

See now that we are also concerned with the angular dependence

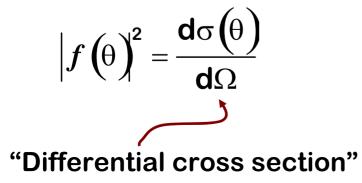
Called "differential cross section"

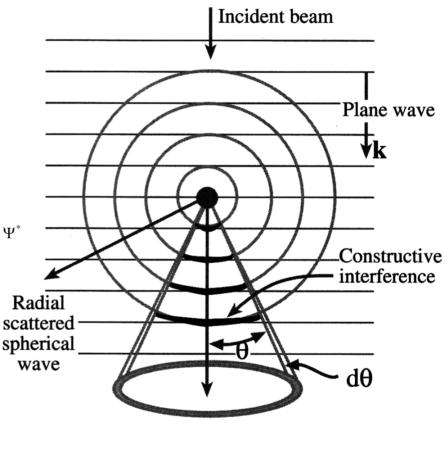


$f(\theta)$ is measure of amplitude of electron wave scattered from the atom

Is thus the tendency of the scattered wave to interfere in a constructive manner I=Ψ×Ψ* with the incident wave (more on this later)

f²(θ) is proportional to the scattered intensity[†]





[†] I= $\Psi \times \Psi^*$

Differential cross sections

The area offered by the scatter (d σ) for scattering the incident electron into a particular increment of solid angle, d Ω

Differential cross sections can be found by solving Schrödinger Eqn inside the atom (!)

There are three primary <u>models</u> used to do this to find differential cross sections:

- "Screened Columbic"
- Thomas-Fermi / Rutherford
- Mott

Lots of heavier physics in this, which I'll ignore

- See both Reimer and Fultz & Howe texts
- Ask me if you are interested, and we can discuss one on one

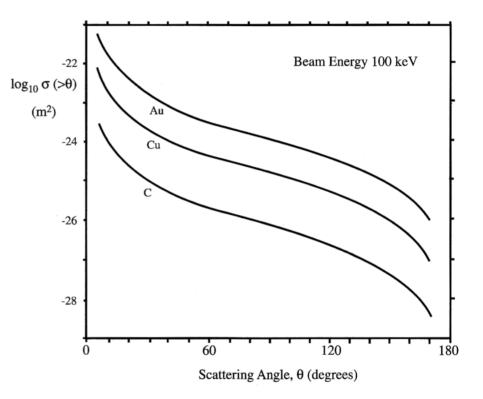
Differential cross section

Has general form shown to the right

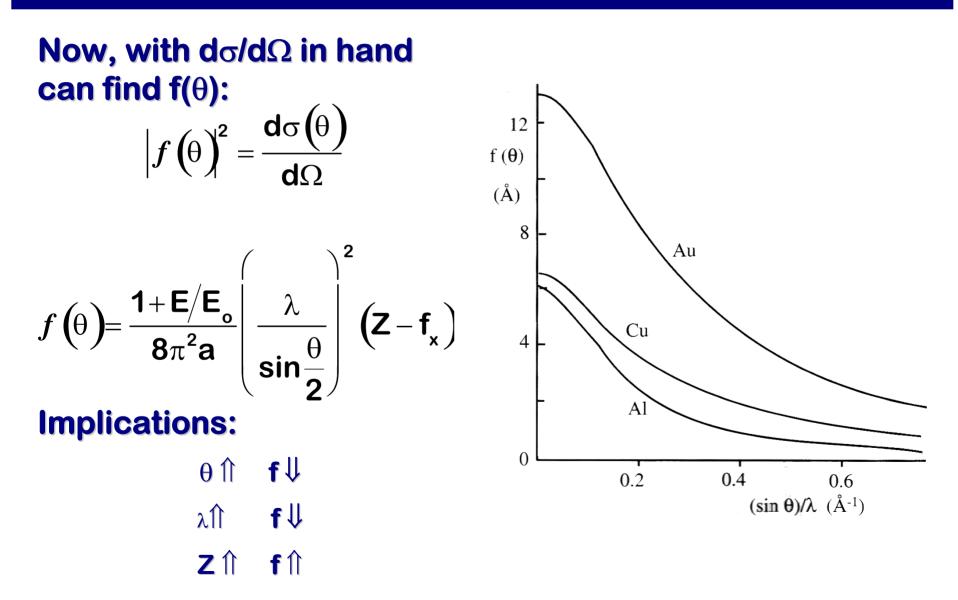
For lower Z & lower V use the "Screened Relativistic Rutherford Cross Section"

$$\frac{d\sigma(\theta)}{d\Omega} = \frac{\lambda_{R}^{4} Z^{2}}{64\pi^{4} (a_{o})^{2} \left(sin^{2} \frac{\theta}{2} + \left(\frac{\theta}{2} \right)^{2} \right)^{2}}$$

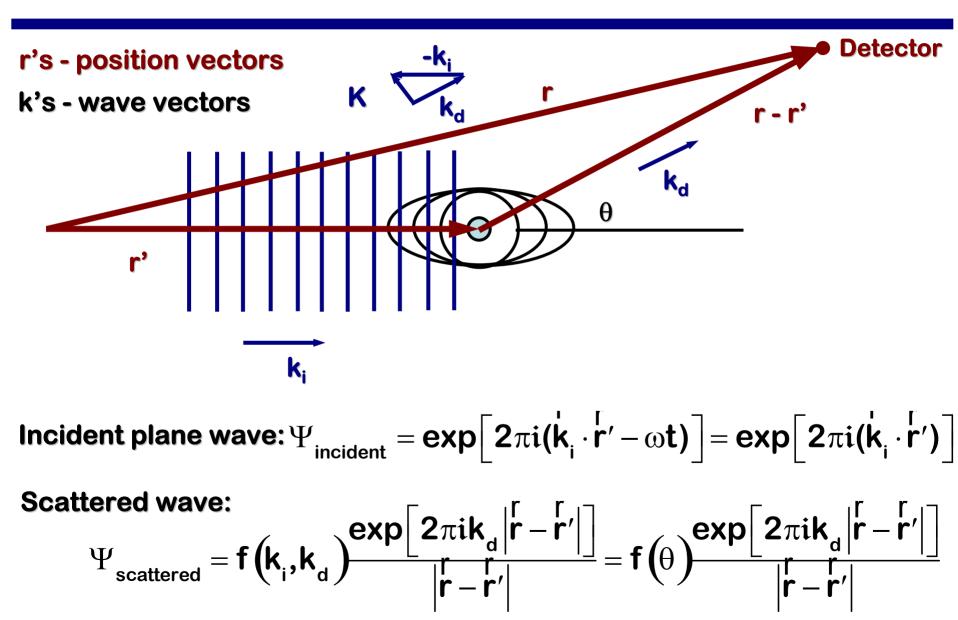
For heavier elements & larger voltages the "Mott Cross Section" is more accurate



Screened Relativistic Rutherford Cross Section vs. Angle



Coherent elastic scattering



Additionally, can use a Green's function approach to formally solve the Schrödinger Eqn. for the electron when it is under the influence of the atomic potential

$$\psi(\mathbf{r}) = \psi_{inc} + \psi_{scatt}$$

= $\exp(2\pi i \mathbf{k}_{i} \cdot \mathbf{r}) + \frac{2m}{h^{2}} \int V(\mathbf{r}') \psi(\mathbf{r}') \mathcal{G}(\mathbf{r}, \mathbf{r}') d^{3}\mathbf{r}'$

In the limit of weak scattering (called "first Born approximation" or "kinematical theory of diffraction"), this Green's function is: $\vec{r} >> \vec{r}'$

$$\mathbf{G}(\mathbf{r},\mathbf{r}') = \frac{-1}{4\pi} \frac{\exp\left[2\pi i \mathbf{k}_{d} \cdot (\mathbf{r} - \mathbf{r}')\right]}{|\mathbf{r}|}$$

Again: See Fultz & Howe or Reimer for details

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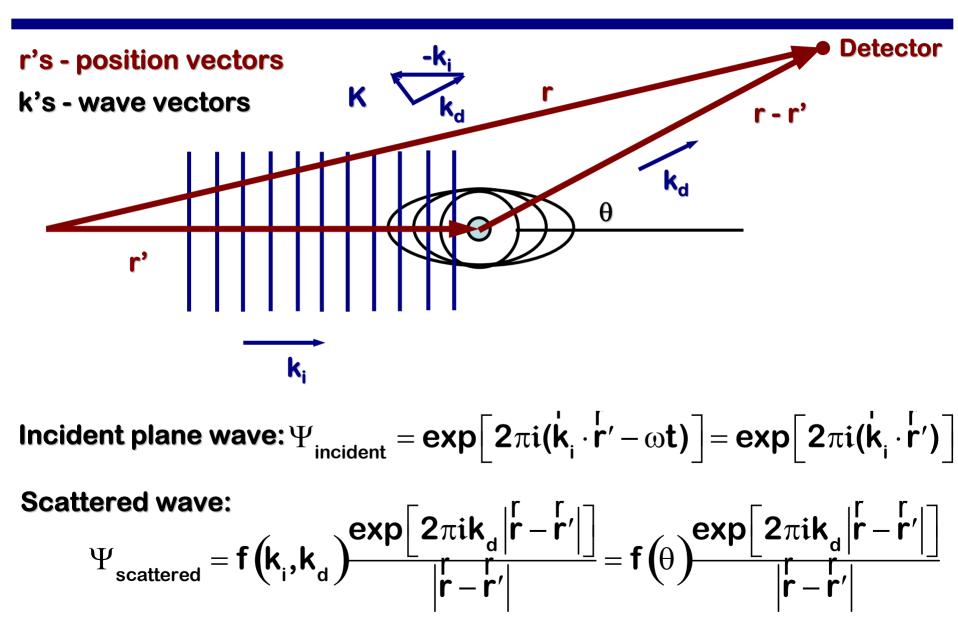
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Substitute:

$$\begin{aligned} \vec{r} \gg \vec{r}' \\
\psi(\vec{r}) &= \psi_{inc} + \psi_{scatt} \\
&= \exp(2\pi i k_{i} \cdot \vec{r}) - \frac{2m}{h^{2}} \int V(\vec{r}') \psi(\vec{r}') G(\vec{r}, \vec{r}') d^{3}\vec{r}' \\
&= \exp(2\pi i k_{i} \cdot \vec{r}') - \frac{-1}{4\pi} \frac{\exp[2\pi i k_{d} \cdot (\vec{r} - \vec{r}')]}{|\vec{r}|} \\
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Atomic form factor

So, with weak scattering, we have:

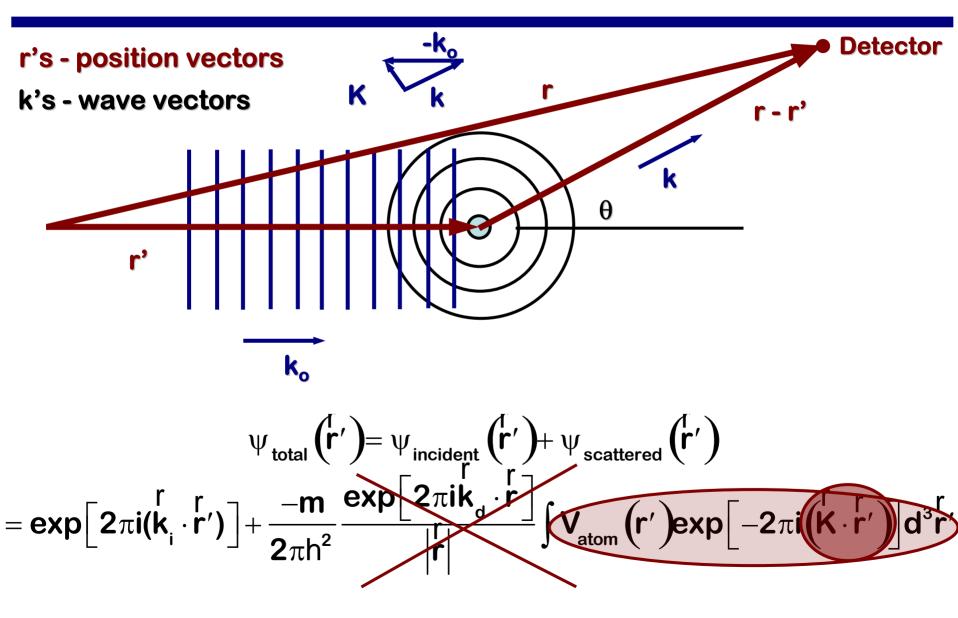
$$\psi_{\text{scatt}} = f_{\text{el}} \begin{pmatrix} r \\ K \end{pmatrix} \frac{\exp(2\pi i \dot{k}_{\text{d}} \cdot \dot{r})}{|\dot{r}|} \qquad \dot{K} = \dot{k}_{\text{i}} - \dot{k}_{\text{d}}$$

$$\psi_{\text{scatt}} = f_{\text{el}} \begin{pmatrix} r \\ K \end{pmatrix} = \frac{-m}{2\pi h^2} \int V_{\text{atom}} (r') \exp\left[-2\pi i \left(\begin{matrix} r \\ K \cdot \dot{r} \end{matrix} \right) \right] d^3 \dot{r}'$$

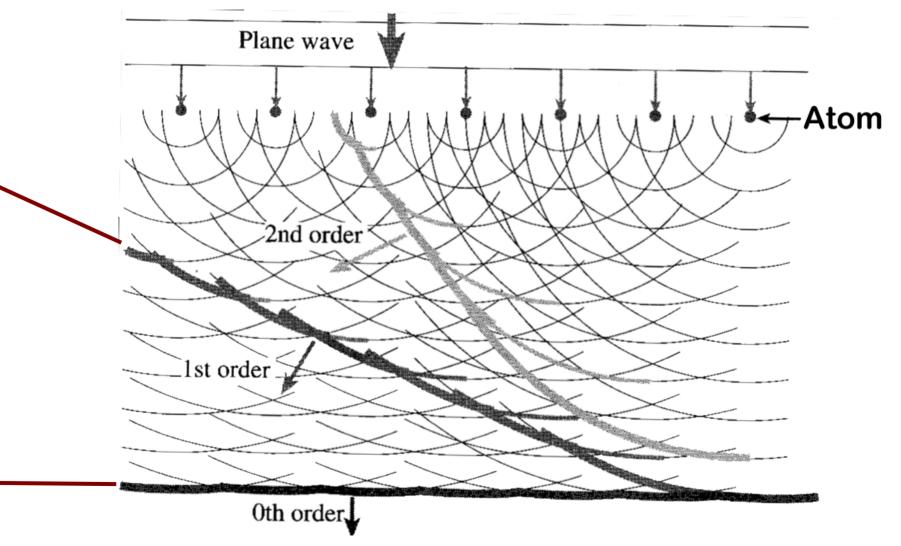
Scattered wave is proportional to Fourier transform of scattering potential

$$\Psi_{\text{scattered}}\left(\vec{\mathbf{K}},\vec{\mathbf{r}}\right) = \frac{-\mathbf{m}}{2\pi\hbar^2} \frac{\exp\left[2\pi i\left(\vec{\mathbf{k}}_{d}\cdot\vec{\mathbf{r}}\right)\right]}{\left|\vec{\mathbf{r}}\right|} \int V_{\text{atom}}\left(\vec{\mathbf{r}}'\right) \exp\left[-2\pi i\left(\vec{\mathbf{K}}\cdot\vec{\mathbf{r}}'\right)\right] d^{3}\vec{\mathbf{r}}'$$

Total wave

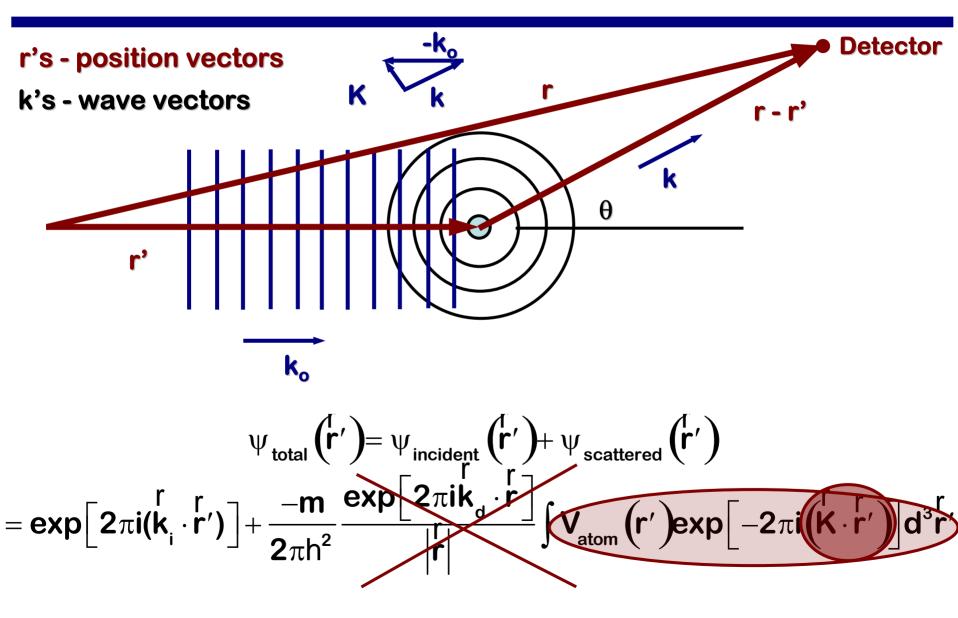


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